

# Shielded Balanced and Coaxial <br> Transmission Lines — Parametric <br> Measurements and Instrumentation <br> Relevant to Signal Waveform Transmission in Digital Service 

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# Shielded Balanced and Coaxial Transmission Lines Parametric Measurements and Instrumentation Relevant to Signal Waveform Transmission in Digital Service 

W. L. Gans
N. S. Nahman

Electromagnetic Technology Division
National Engineering Laboratory
National Bureau of Standards
Boulder, Colorado 80303
U.S. DEPARTMENT OF COMMERCE, Malcolm Baldrige, Secretary

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# SHIELDED BALANCED AND COAXIAL TRANSMISSION LINES PARAMETRIC MEASUREMENTS AND INSTRUMENTATION RELEVANT TO SIGNAL WAVEFORM TRANSMISSION IN DIGITAL SERVICE 

W. L. Gans and N. S. Nahman


#### Abstract

A method is presented for determining the impulse and step responses of a shielded cable using time domain terminal measurements and a physically based mathematical model for the transmission line properties of the cable. The method requires a computer controlled time domain measurement system and was implemented using the NBS Automatic Pulse Measurement System (APMS). Data are also developed for the frequency domain complex propagation function (attenuation and its related minimum-phase shift). The method is applied to 12 shielded paired-conductor (balanced) cables and 5 coaxial cables. Time domain responses are presented for three nominal cable lengths, 60 m ( 200 ft ), 150 m ( 500 ft ), and 300 m ( 1000 ft ). The time domain responses are applied to the estimation of bit error rate increases due to the insertion of the cables into a digital signaling system employing a balanced polar NRZ waveform. Also discussed is the application of the time domain responses to time domain reflectometry techniques for cable acceptance tests and field-site testing of installed cables.


Key words: bit error rate; coaxial transmission lines; digital communication impulse response; instrumentation; measurements; shielded balanced transmission lines; time domain; transmission line model; waveform.

## 1. INTRODUCTION

### 1.1. Background

Nominal frequency response data supplied by cable manufacturers are not sufficient for the complete evaluation or simulation of cables (transmission lines) in digital systems. Data which must be acquired prior to engineering analysis for digital transmission are the driving-point and transfer impulse responses. These are most easily obtained through pulse or time domain testing.

Furthermore, time domain measurement methods and performance data for the characterization of balanced digital-service, shielded transmission lines [ 60 m ( 200 ft ) to 300 m ( 1000 ft ) lengths] are required for the design, signal analysis, and circuit monitoring relevant to the digital-service transmission line interconnections between the Technical Control Facility (TCF) and the Radio/Microwave Exchange (MUX) as encountered in U.S. Army telecommuncation installations. Also, time domain measurement methods are required for the identification and location of installation and/or operating faults which can impair the digital waveform fidelity, which in turn can cause a degradation of the digital signaling capacity. Such measurement methods and their interpretation fundamentally depend upon a knowledge of typical impulse responses.

Because response data suitable for digital cable simulation was not available, the U.S. Army Communications - Electronics Engineering Installation Agency (USACEEIA) sought assistance from the National Bureau of Standards (NBS) to develop the required data for a specified group of commercially available cables. The NBS Time Domain Metrology Group was uniquely qualified to render such assistance
by virtue of their experience in (1) the development and use of the NBS Automatic Pulse Measurement System and (2) the characterization of transient responses for shielded cables. The subject report summarizes the work performed by the NBS Time Domain Metrology Group for USACEEIA and is for the most part a catalog of cable impulse responses and associated transfer functions.

### 1.2 Project Objective

The overall project objective contained three components; listed in order of their relative technical priority, they were:

1. Determine and compile from measurements the parametric data necessary for engineering design and computer simulation of typical shielded cables for digital service. Demonstrate the use of the data to provide time domain and frequency domain simulation.
2. Simulate time domain reflectometry reference signatures suitable for cable field testing and fault location.
3. Demonstrate a time domain data acquisition system capable of (a) acquiring the signal waveforms of a transmission line under test conditions and (b) provide measured cable parameters suitable for subsequent batch processing in a central computer system for engineering design and simulation studies.

Originally, it was proposed to accomplish each one of the above components individually in a specific phase of the work during the contract duration. However, to allow for cable and component delivery delays, and the natural interaction between the objective components, the work was not rigidly accomplished in specific phases. The subject report presents the results of the work in accomplishing objectivecomponents (1), (2), and (3) above.

### 1.3 Report Organization

This report is divided into eight chapters and two appendices, the first chapter being the present introduction. Chapter 2 presents the transmission line model, the definition of parameters, and a discussion of the model frequency dependence.

Chapter 3 describes the measurement methods used to obtain the frequency and time domain data required for the calculation of the transmission line parameters.

Chapter 4 provides a qualitative overview of the computer calculations for (1) the extraction or determination of the transmission line parameters, (2) the impulse response, and (3) the step response.

Chapter 5 presents an example of the typical results for the determination of the transmission line parameters and the corresponding frequency and time domain responses for a given cable.

Chapter 6 describes a method of estimating the bit-error-rate degradation of a balanced polar NRZ waveform produced by a given length of cable.

Chapter 7 discusses the time domain reflectometry method for testing cables which is suitable for acceptance and field-site tests. Typical results are given for the worst case situations of open and short-circuit discontinuities.

Appendix A is a catalog of complex phasor transfer functions, their associated time domafn responses (both reflection and transmission) for selected commercially available shielded cables of different lengths, usually $60 \mathrm{~m}(200 \mathrm{ft}), 150 \mathrm{~m}(500 \mathrm{ft})$, and $300 \mathrm{~m}(1000 \mathrm{ft})$. The cables included are:
A. Shielded Paired-Conductor Transmission Lines: RG 22, U.S. Army $W D-37$, and ten commercially available cables designated as cables A through J.
B. Coaxial-conductor Transmission Lines: RG58C, RG59B, RG214, RG223, and a commercially available triaxial cable designated as cable K.

Appendix B contains the computer programs which were used in the NBS APMS to accomplish the cable measurements, characterization and simulation. Appendix $C$ is a description of one design for a time domain data acquisition system (NBS APMS).

## 2. DESCRIPTION OF THE TRANSMISSION LINE MODEL AND THE DEFINITION OF THE MODEL PARAMETERS

In this chapter a brief discussion of the general transmission line equations is presented so as to provide a basis for the formulation of the transmission line simulation model. Following the introductory discussion, the simulation model parameters are defined and the transmission line equations are expressed in terms of the model parameters. Finally, various limiting cases of the model versus frequency are discussed to provide insight as to the correlation of the model with experiment.

### 2.1 The General Transmission Line Equations

A uniform transmission line can be characterized in terms of a series impedance per unit length, $Z(s)$, and a shunt admittance per unit length, $Y(s), f i g$. 2.1. This follows from the uniformity property which simply means that the transmission line geometric form and material composition does not change with length, i.e., any segment of the transmission line possesses the same values of $Z(s)$ and $Y(s)$. The specific functional forms of $Z(s)$ and $Y(s)$ depend upon geometric form and material parameters. In practice $Z(s)$ and $Y(s)$ are combined to form two other functions which naturally appear in the transmission line equations and are called the characteristic impedance $Z_{0}(s)$ and the propagation function $\gamma(s)$.

$$
\begin{array}{ll}
Z_{o}(s) & =\sqrt{\frac{Z(s)}{Y(s)}} \quad, \text { ohms } \\
Y(s) & =\sqrt{Z(s) Y(s)} \tag{2-2}
\end{array}, \text { per unit length } .
$$

For a transmission line of length $\ell$ connected between a generator $E_{g}(s)$ of impedance $Z_{g}(s)$ and $a$ load impedance $Z_{\ell}(s)$, fig. 2.2, the transfer and driving point equations are given by [1], respectively,

$$
\begin{equation*}
\frac{E_{r}(s)}{E_{g}(s)}=2\left[\frac{Z_{0}(s)}{Z_{g}(s)+Z_{o}(s)}\right]\left[\frac{Z_{\ell}(s)}{Z_{\ell}(s)+Z_{o}(s)}\right]\left[\frac{e^{-\ell \gamma(s)}}{1-\rho_{s}(s) \rho_{r}(s) e^{-2 \ell \gamma(s)}}\right] \tag{2-3}
\end{equation*}
$$

$$
\begin{equation*}
\frac{E_{S}(s)}{I_{S}(s)}=Z_{o}(s) \quad \frac{1+\rho_{r}(s) e^{-2 \ell \gamma(s)}}{1-\rho_{r}(s) e^{-2 \ell \gamma(s)}} \tag{2-4}
\end{equation*}
$$

where the reflection coefficients $\rho_{S}(s) \rho_{r}(s)$ are obtained from

$$
\begin{equation*}
\rho_{i}(s)=\frac{Z_{i}(s)-Z_{o}(s)}{Z_{i}(s)+Z_{o}(s)} \tag{2-5}
\end{equation*}
$$

These equations describe a transmission line network which is not terminated at either end in the transmission line characteristic impedance.

When a termination $Z_{i}(s)$ is equal to $Z_{o}(s)$, then the corresponding $\rho_{i}(s)$ is zero, and the transmission line is said to be matched at that end. There are three possible situations for matched terminations:

1. Doubly matched $Z_{g}(s)=Z_{\ell}(s)=Z_{o}(s)$

$$
\begin{gather*}
\left.\frac{E_{r}(s)}{E_{g}(s)}\right]=\frac{e^{-2 \gamma(s)}}{2}  \tag{2-6}\\
\rho_{g}=\rho_{r}=0
\end{gather*}
$$

$$
\begin{array}{r}
\left.\frac{E_{s}(s)}{I_{s}(s)}\right]=Z_{o}(s)  \tag{2-7}\\
\rho_{g}=\rho_{r}=0
\end{array}
$$

2. Matched Load $Z_{\ell}(s)=Z_{o}(s)$

$$
\begin{equation*}
\left.\frac{E_{r}(s)}{E_{g}(s)}\right]_{\rho_{\ell}=0}=\left[\frac{Z_{o}(s)}{Z_{g}(s)+Z_{o}(s)}\right] e^{-\ell \gamma(s)} \tag{2-8}
\end{equation*}
$$

$$
\begin{gather*}
\left.\frac{E_{s}(s)}{I_{s}(s)}\right]=Z_{0}(s)  \tag{2-9}\\
\rho_{\ell}=0
\end{gather*}
$$

3. Matched Generator $Z_{g}(s)=Z_{o}(s)$

$$
\begin{align*}
& \left.\frac{E_{r}(s)}{E_{g}(s)}\right]=\left[\frac{Z_{\ell}(s)}{Z_{\ell}(s)+Z_{o}(s)}\right] e^{-\ell \gamma(s)}  \tag{2-10}\\
& \left.\frac{E_{g}(s)}{I_{s}(s)}\right]_{\rho_{g}}=0 \tag{2-11}
\end{align*}
$$

Equations (2-3) through (2-11) reduce to the lossless or ideal cases when

$$
\begin{equation*}
\mathrm{Z}(\mathrm{~s})=\mathrm{sL} \tag{2-12}
\end{equation*}
$$

and

$$
\begin{equation*}
Y(s)=s C \tag{2-13}
\end{equation*}
$$

which in turn yield a delay propagation function and a resistive characteristic impedance, respectively,

$$
\begin{equation*}
\gamma(s)=s \sqrt{L C} \tag{2-14}
\end{equation*}
$$

and

$$
\begin{equation*}
Z_{0}(s)=R_{0}=\sqrt{\frac{L}{C}} \tag{2-15}
\end{equation*}
$$

Under such conditions, there is no distortion with propagation, and also the transmission line can be terminated with a resistor, $R_{o}$, for reflection-free transmission.

For practical shielded cables, there are losses present, which in turn, means that both $\gamma(s)$ and $Z_{o}(s)$ will be complicated functions of $s$. The resultant propagation along the line will distort pulses due to the nature of $\gamma(s)$. Also, because $Z_{0}(s)$ is $s$ dependent, resistive terminations $R_{g}$, $R_{2}$ will not match (equal) the characteristic impedance; under such conditions in equations (2-3) through (2-11), the voltage divider terms and their combinations which form the reflection coefficients,

$$
\frac{Z_{0}(s)}{R_{g}+Z_{o}(s)} ; \quad \frac{Z_{\ell}(s)}{R_{\ell}+Z_{o}(s)} ; \quad \frac{R_{i}-Z_{o}(s)}{R_{i}+Z_{0}(s)}
$$

contribute to the overall distortion of the transmitted pulse as they are not independent of $s$.

### 2.2 The $s^{\mathrm{m}}$ Model and Its Parameters

The physically-based mathematical model used in this work was originally developed by Nahman in 1962 for the sub-nanosecond to nanosecond domain transient analysis of coaxial-conductor cables [2]. Subsequent investigation by Nahman and Riad using the NBS APMS showed that the model could also be used to simulate shielded paired-conductor cables of parallel or twisted conductor construction [3].

The model specifies $Z(s)$ and $Y(s)$ as

$$
\begin{equation*}
\mathrm{Z}(\mathrm{~s})=\mathrm{R}+\mathrm{sL}+K \mathrm{~s}^{\mathrm{m}}, \quad 0<\mathrm{m}<1 \tag{2-16}
\end{equation*}
$$

and

$$
\begin{equation*}
Y(s)=s C \tag{2-17}
\end{equation*}
$$

where
$R \equiv$ dc resistance/unit length
$L \equiv$ Inductance/unit length
$K \equiv$ High frequency loss coefficient/unit length
$m \equiv$ High frequency loss exponent
$C \equiv$ Capacitance/unit length

Consequently, the model equivalent circuit per unit length is as shown in fig. 2.3. The propagation function and the characteristic impedance are then given by
and

$$
\begin{equation*}
\gamma(s)=\left[\left(R+s L+K s^{m}\right) s C\right]^{1 / 2}, \quad 0<m<1 \tag{2-23}
\end{equation*}
$$

$$
\begin{equation*}
Z_{o}(s)=\left[\left(R+s L+K s^{m}\right) / s C\right]^{1 / 2}, \quad 0<m<1 \tag{2-24}
\end{equation*}
$$

In the high frequency region, i.e., when $s$ is such that

$$
\begin{equation*}
\text { sL >> } \mathrm{Ks}^{\text {m }} \gg \mathrm{R} \tag{2-25}
\end{equation*}
$$

the model series impedance is closely approximated by

$$
\begin{equation*}
\underset{H F}{Z(s)]}=s L+K s^{m} \tag{2-26}
\end{equation*}
$$

which in turn gives for $\gamma(s)$ and $Z_{o}(s)$

$$
\begin{equation*}
\gamma\left(\underset{H F}{j}=s \sqrt{L C}\left(1+\frac{K^{m}}{s L}\right)^{1 / 2}\right. \tag{2-27}
\end{equation*}
$$

and

$$
\begin{equation*}
\left.\mathrm{Z}_{\mathrm{o}}(\mathrm{~s})\right]=\sqrt{\frac{\mathrm{L}}{\mathrm{C}}}\left(1+\frac{\mathrm{Ks}^{\mathrm{m}}}{\mathrm{sL}}\right)^{1 / 2} . \tag{2-28}
\end{equation*}
$$

Because

$$
\begin{equation*}
\frac{\mathrm{Ks}^{\mathrm{m}}}{\mathrm{sL}} \ll 1 \tag{2-29}
\end{equation*}
$$

then

$$
\begin{equation*}
\left(1+\frac{\mathrm{Ks}^{\mathrm{m}}}{\mathrm{sL}}\right)^{1 / 2} \doteq 1+\frac{\mathrm{Ks}^{\mathrm{m}}}{2 \mathrm{sL}} \tag{2-30}
\end{equation*}
$$

Defining

$$
\begin{equation*}
\delta(s)=\frac{\mathrm{Ks}^{\mathrm{m}}}{2 \mathrm{sL}} \tag{2-31}
\end{equation*}
$$

then the high frequency $\gamma(s)$ and $Z_{o}(s)$ may be expressed as

$$
\begin{equation*}
\gamma(s)]_{H F}=s \sqrt{L C}[1+\delta(s)],|\delta(s)| \ll 1 \tag{2-32}
\end{equation*}
$$

and

$$
\begin{equation*}
\left.Z_{o}(s)\right]_{\mathrm{HF}}=\sqrt{\frac{\mathrm{L}}{\mathrm{C}}}[1+\delta(\mathrm{s})],|\delta(\mathrm{s})| \ll 1 . \tag{2-33}
\end{equation*}
$$

In the limit as $|s|$ approaches infinity, (2-32) and (2-33) yield the time delay per unit length and the high frequency or geometric or nominal characteristic impedance, respectively,

$$
\begin{align*}
& \lim _{|s| \rightarrow \infty}[\gamma(s)]_{H F}=s \sqrt{L C}  \tag{2-34}\\
& \lim _{|s| \rightarrow \infty}\left[Z_{o}(s)\right]_{H F}=\sqrt{\frac{L}{C}} \tag{2-35}
\end{align*}
$$

which defines the time delay per unit length

$$
\begin{equation*}
T=\sqrt{L C} \tag{2-36}
\end{equation*}
$$

and the high frequency characteristic impedance,

$$
\begin{equation*}
R_{o}=\sqrt{\frac{L}{C^{\prime}}} \tag{2-37}
\end{equation*}
$$

## 2.3 s-Domain Circuit Properties of the Model

Consider the circuit of fig. 2.4 which is the circuit of fig. 2.2 with

$$
\begin{align*}
& Z_{g}(s)=Z_{\ell}(s)=R_{o}  \tag{2-38}\\
& \gamma(s)=\left[\left(R+s L+K s^{m}\right) s c\right]^{1 / 2}  \tag{2-23}\\
& Z_{o}(s)=\left[\left(R+s L+K s^{m}\right) / s c\right]^{1 / 2} \tag{2-24}
\end{align*}
$$

Here, the transmission line is terminated at each end in the nominal or high frequency characteristic impedance, $R_{o}$. This is a doubly mismatched situation for low frequencies, but this is the usual case encountered in practice. In all that follows in this report, the generator and load impedances will be equal to $R_{0}$. Consequently, equations (2-3) through (2-5) become

$$
\begin{align*}
& \frac{E_{r}(s)}{E_{g}(s)}=2\left[\frac{Z_{0}(s)}{R_{0}+Z_{0}(s)}\right]\left[\frac{R_{0}}{R_{0}+Z_{o}(s)}\right]\left[\frac{e^{-\ell \gamma(s)}}{1-[\rho(s)]^{2} e^{-2 \ell \gamma(s)}}\right]  \tag{2-39}\\
& \frac{E_{S}(s)}{I_{s}(s)}=Z_{0}(s) \frac{1+\rho(s) e^{-2 \ell \gamma(s)}}{1-\rho(s) e^{-2 \ell \gamma(s)}}  \tag{2-40}\\
& \rho_{S}(s)=\rho_{r}(s)=\rho(s)=\frac{R_{0}-Z_{0}(s)}{R_{0}+Z_{o}(s)} \tag{2-41}
\end{align*}
$$

For the purposes of the following discussion, it is convenient to express equations (2-38) and (2-39) in terms of the hyperbolic functions

$$
\begin{align*}
& \cosh \ell_{\gamma}(s)=\frac{e^{\ell \gamma(s)}+e^{-\ell \gamma(s)}}{2}  \tag{2-42}\\
& \sinh \ell_{\gamma}(s)=\frac{e^{\ell \gamma(s)}-e^{-\ell \gamma(s)}}{2} \tag{2-43}
\end{align*}
$$

Doing so, there results

$$
\begin{equation*}
\frac{E_{r}(s)}{E_{g}(s)}=\frac{2 Z_{o}(s) R_{0}}{2\left[R_{0}^{2}+Z_{o}^{2}(s)\right] \sinh \operatorname{lr}(s)+4 R_{o}} \frac{Z_{o}(s) \cosh \operatorname{lr}(s)}{} \tag{2-44}
\end{equation*}
$$

$$
\begin{equation*}
\frac{E_{s}(s)}{I_{s}(s)}=Z_{o}(s) \frac{R_{0} \cosh \ell \gamma(s)+Z_{0}(s) \sinh \ell \gamma(s)}{Z_{0}(s) \cosh \ell \gamma(s)+R_{0} \sinh \ell \gamma(s)} \tag{2-45}
\end{equation*}
$$

Consideration of fig. 2.4 for $s=0$ (dc) leads to the network shown in fig. 2.5 and the following conclusions:

1. The transfer function $E_{r}(o) / E_{g}(o)$, is equal to $R_{o} /\left(2 R_{o}+\ell R\right)$.
2. The impedance looking into the line, $E_{S}(0) / I_{S}(0)$, is equal to $\ell R+R_{0}$.

To demonstrate that these two conclusions are consistent with the mathematical model, consider (2-44) and (2-45) when $|s| \rightarrow o$. First of all,

$$
\left.\begin{array}{l}
Z_{0}(s) \rightarrow \sqrt{\frac{R}{s C}}  \tag{2-46}\\
|s| \rightarrow 0
\end{array}\right\} \begin{aligned}
& \substack{s,|s| \rightarrow 0}
\end{aligned}
$$

$\underset{|s| \rightarrow 0}{\lim Z_{o}(s) \rightarrow \infty}$

Putting (2-46) and (2-49) into (2-44) and (2-45) gives for $|s| \rightarrow 0$

$$
\begin{equation*}
\left[\frac{E_{r}(s)}{E_{g}(s)}\right]_{|s| \rightarrow 0} \rightarrow \frac{2 R_{o} \sqrt{s R C}}{2\left(R_{o}^{2} s C+R\right) \sinh \ell \sqrt{s c R}+4 R_{o} \sqrt{s c R} \cosh \ell \sqrt{s c R}} \tag{2-50}
\end{equation*}
$$

$$
\begin{equation*}
\left[\frac{E_{s}(s)}{I_{s}(s)}\right]_{|s|_{\rightarrow 0}} \rightarrow \frac{\sqrt{s C R} R_{0} \cosh \ell \sqrt{s c R}+R \sinh \ell \sqrt{s c R}}{\sqrt{s c R} \cosh \ell \sqrt{s c R}+R_{0} s c \sinh \ell \sqrt{s c K}} . \tag{2-51}
\end{equation*}
$$

When $|s|=0,(2-50)$ and (2-51) both yield $o / o$ and are indeterminate. After differentiating the numerators and the denominators, and again applying the limit, the results are

$$
\begin{align*}
& \lim _{|s| \rightarrow 0}\left[\frac{E_{r}(s)}{E_{g}(s)}\right]=\frac{R_{0}}{2 R_{0}+\ell R}  \tag{2-52}\\
& \left\lvert\, \lim _{\mid \rightarrow 0}\left[\frac{E_{s}(s)}{I_{s}(s)}\right]=R_{0}+\ell R\right. \tag{2-53}
\end{align*}
$$

which agrees with the earlier conclusions.
Next, consider the high frequency behavior of the circuit. Putting (2-32) and (2-33) into (2-44) and (2-45) yie1ds

$$
\begin{align*}
\left.\frac{E_{r}(s)}{E_{g}(s)}\right] & =\frac{1}{e^{s T l}[1+\delta(s)]}+\frac{[\delta(s)]^{2}}{2[1+\delta(s)]} \sinh \{\operatorname{sTl}[1+\delta(s)]\} \\
& =e^{-s T l[1+\delta(s)]} \tag{2-54}
\end{align*}
$$

and

$$
\begin{align*}
\left.Z_{S}(s)\right] & =R_{0}[1+\delta(s)] \frac{e^{s T l[1+\delta(s)]}+\delta(s) \sinh \{s T \ell[1+\delta(s)]\}}{e^{s T \ell[1+\delta(s)]}+\delta(s) \cosh \{s T \ell[1+\delta(s)]\}} \\
& =R_{0}[1+\delta(s)] \tag{2-55}
\end{align*}
$$

as $|\delta(s)| \ll 1$. Note that $Z_{s}(s)$ is equal to $\left.Z_{o}(s)\right]$.

### 2.4 High Frequency Attenuation and Phase Shift

Equation (2-54) can be written as

$$
\begin{equation*}
\frac{E_{r}(s)}{E_{g}(s)}=e^{-s T \ell} e^{-s T \ell \delta(s)} \tag{2-56}
\end{equation*}
$$

where the first factor represents the time delay of the cable (linear phase shift in the frequency domain), while the second contains the distortion due to high frequency losses. The inverse Laplace transform of (2-56) would yield the initial or early time response of the circuit impulse response. Consequently, (2-56) exhibits the time delay and then the initial impulse response.

In the frequency domain the transfer function is

$$
\frac{E_{r}(j \omega)}{E_{g}(j \omega)}=e^{-j \omega \ell T} e^{\frac{-\ell T K}{2 L}(j \omega)^{m}}
$$

or from (2-36) and (2-37),

$$
\begin{equation*}
\frac{E_{r}(j \omega)}{E_{g}(j \omega)}=e^{-j \omega \ell T} e^{-\frac{\ell K}{2 R_{o}}(j \omega)^{m}} \tag{2-57}
\end{equation*}
$$

$$
o<m<1
$$

The attenuation $\alpha_{1}(\omega)$ and the phase $\beta(\omega)$ of the transfer function are defined by

$$
\begin{equation*}
\frac{E_{r}(j \omega)}{E_{g}(j \omega)}=e^{-\alpha}(\omega)-j \beta(\omega) \tag{2-58}
\end{equation*}
$$

and are

$$
\begin{align*}
& \alpha_{1}(\omega)=\text { Real Part of } \frac{\ell K}{2 R_{o}}(j \omega)^{m}  \tag{2-60}\\
& B(\omega)=\omega \ell T+\text { Imaginary Part of } \frac{\ell K}{2 R_{o}}(j \omega)^{m} . \tag{2-61}
\end{align*}
$$

In the phase $\beta(\omega)$, the first term is the linear phase shift associated with the transmission line time delay. In network terminology $\exp (-j \omega \ell T)$ is an all-pass function possessing a linear phase shift. The second term in the phase is uniquely related to the attenuation because the two terms are the real and imaginary parts of an analytic function along the frequency axis, $j \omega$, of the complex frequency plane $s=\sigma+j \omega$; thus the factor in (2-57)

$$
\begin{equation*}
F_{1}(j \omega)=e^{-\frac{\ell K}{2 R_{o}}(j \omega)^{m}} \quad 0<m<1 \tag{2-61}
\end{equation*}
$$

is a minimum phase function. Hence, it is seen that the frequency dependent attenuation has associated with it a frequency dependent (non-linear) phase function.

Defining

$$
\begin{equation*}
\beta(\omega)=\beta_{0}(\omega)+\beta_{1}(\omega) \tag{2-62}
\end{equation*}
$$

where

$$
\begin{align*}
& \beta_{0}(\omega)=\omega \ell T  \tag{2-63}\\
& \beta_{1}(\omega)=\text { Imaginary Part of } \frac{-\ell K}{2 R_{o}}(j \omega)^{m} \tag{2-64}
\end{align*}
$$

the minimum phase complex exponent is then

$$
\begin{equation*}
A(j \omega)=\alpha_{1}(\omega)+j \beta_{1}(\omega) \tag{2-65}
\end{equation*}
$$

It is equal to the exponent of (2-57) with the linear phase shift removed, i.e.,

$$
\begin{align*}
& \ell n\left\{e^{-j \omega \ell T} \frac{E_{g}(j \omega)}{E_{r}(j \omega)}\right\}=\alpha_{1}(\omega)+j \beta_{1}(\omega) \\
& \alpha_{1}(\omega)+j \beta_{1}(\omega)=\frac{\ell K}{2 R_{o}}(j \omega)^{m} \tag{2-66}
\end{align*}
$$

$$
\begin{equation*}
\alpha_{1}(\omega)+j \beta_{1}(\omega)=\frac{\ell K}{2 R_{o}} \quad \omega^{m} \quad\left(\cos \frac{m \pi}{2}+j \sin \frac{m \pi}{2}\right) \tag{2-67}
\end{equation*}
$$

.Hence,

$$
\begin{equation*}
\alpha_{1}(\omega)=\frac{\ell K}{2 R_{o}} \omega^{m} \cos m \pi / 2 \text {, nepers/unit length } \tag{2-68}
\end{equation*}
$$

and

$$
\begin{equation*}
\beta_{1}(\omega)=\frac{\ell K}{2 R_{o}} \quad \omega^{m} \text { sin } m \pi / 2 \text {, radians/unit length } \tag{2-69}
\end{equation*}
$$

which shows that both the attenuation and phase vary as $\omega^{m}$. Their ratio is

$$
\begin{equation*}
\frac{\alpha_{1}(\omega)}{\beta_{1}(\omega)}=\cot m \pi / 2 \tag{2-70}
\end{equation*}
$$

Figure 2.6 shows a log-log graph illustrating equations (2-68) through (2-70). Note that the linear slope (i.e., the measured slope using a linear scale) is equal to the fractional power, m. Similarly, a graph of the attenuation $\alpha_{1}(\omega)$ in $d B$ ( 8.68 nepers equals one $d B$ ) or the minimum phase $\beta_{1}(\omega)$ in degrees ( 0.0175 radians equals one degree) would have a linear slope of m .

For frequencies ranging from $D C$ to very high frequencies, the log-log attenuation graph for the model has two asymptotes, fig. 2.7. For very low frequencies the attenuation asymptotically approaches the value set by the transfer ratio (2-52), i.e.,

$$
\begin{equation*}
\alpha(o)=20 \log \frac{2 R_{o}+2 R}{R_{o}} \tag{2-71}
\end{equation*}
$$

On the other hand, for very high frequencies $\alpha(\omega)$ asymptotically approaches the fractional power response (2-68); in $d B$, the asymptotic response is

$$
\begin{equation*}
\left.\alpha_{1}(\omega)\right]_{d B}=8.68 \frac{\ell K}{2 R_{o}}|\omega|^{m} \cos m \pi / 2, d B \tag{2-72}
\end{equation*}
$$

Consequently, the model high frequency loss parameter $K$ and the exponent m can be determined from experimental data in the asymptotic high frequency region, provided, of course, that the model is consistent with the physical situation.

In practice, some care must be taken in fitting the model to the experimental data. Specifically, the slope of log-log plot in the high frequency region may not be constant, fig. 2.8 ; this in turn requires a judgement as to how to select a suitable constant slope curve which is tangent to the log-log attenuation curve in the high frequency region, fig. 2.8. Because the NBS measurement system has a dynamic range of better than 40 dB , the experimental data in the 0 to 30 dB range will possess a good signal to noise ratio. Consequently, the determination of $m$ and $K$ is accomplished by fitting in the 30 dB range. The effectiveness of the fit can only be evaluated by comparing the simulated output waveform based upon the modeled impulse response with that of the observed waveform as will be demonstrated in Chapters 4 and 5.

Finally, one other comment should be made regarding the $s^{m}$ model. The shunt admittance/unit length $Y(s)$ is specified as being a pure susceptance $s C,(2-17)$; thus $Y(s)$ contains no conductance $G$ which would represent one particular form of dielectric loss. However, in the computer program written for this work $Y$ (s) has the provision for the form

$$
\begin{equation*}
Y(s)=G+s C \tag{2-73}
\end{equation*}
$$

so that such a loss form can be easily included if needed. In the cables considered here, $G$ was negligible (as it was infinitesimal, or for all practical purposes, zero). This is not to say that the $s{ }^{m}$ model excludes all dielectric loss effects when $G=0$. To the contrary, the slope $m$ and the coefficient $K$ can be effected by high frequency losses due to small dielectric losses in combination with conductor losses.


Figure 2.1 The transmission line equivalent circuit per unit length.


Figure 2.2 The basic transmission circuit consisting of a generator, a transmission line, and a load.


Figute 2.3 The equivalent circuit per unit length for the physicallybased mathematical model.


Figure 2.4 The transmission line circuit terminated at each end in $\mathrm{R}_{0}$, the nominal characteristic impedance. This is a doubly mismatched transmission line because $Z_{o}(s)$ is not equal to $R_{o}$ for all s.


Figure 2.5 The transmission line circuit when $s=0$ (dc)


Figure 2.6 Log-log graph of the minimum phase attenuation $\alpha_{1}(\omega)$ and phase $\beta_{1}(\omega)$. The slope $y / x$ as measured with a linear scale is equal to m .

gp '(m) o8o
$20 \log \left(\frac{2 R_{o}+\ell R}{R_{o}}\right)$


Figure 2.8 Two constant slope curves tangent to a log-log attenuation vs. (angular) frequency curve.

## 3. MEASUREMENT METHODS FOR THE DETERMINATION OF THE TRANSMISSION LINE PARAMETERS

In this chapter the measurement methods required to determine the transmission line model parameters are presented. A brief discussion is presented on the need for measurements and some relations between the basic parameters. Then specific measurement methods are described for the determination of $R$, $C$, $R_{o}, m$ and $K$. $L$ is determined from $C$ and $R_{0}$.

### 3.1 The Need for Measurements

A commercially available or military cable may be characterized in terms of a set of nominal frequency domain transmission line parameters. Typically, they are

```
                                    R - resistance/unit length
                                    C - capacitance/unit length
                                    * s ..
R
    geometric characteristic impedance).
\mp@subsup{\alpha}{1}{}}(f)-\quad\mathrm{ attenuation frequency, f.
```

These parameters are representative of the particular cable type and have been derived from measurements on many different runs of cable. The values can be significantly different from those of a single cable sample and also can vary from manufacturer to manufacturer. Consequently, for precise estimates of the frequency domain performance of a given length of cable, the parameter values should be experimentally determined, i.e., measured. For system design applications the nominal parameter values are used with the hope that they are representative design center values. The main point being stressed here is that the manufacturer's technical specifications are always representative values; individual samples may vary.

For time domain applications, knowledge of the parameters $R, C, R_{o}$ and $\alpha_{1}$ (f) are also required along with the phase shift, $\beta_{1}(f)$, which is uniquely related to $\alpha_{1}(f)$. $\beta_{1}$ (f) is the minimum phase shift component of the total cable phase shift, $\beta(f),(2-62)$ written explicitly in terms of frequency, $f$, i.e., $f=\omega / 2 \pi$.

$$
\begin{align*}
\beta(f) & =\beta_{o}(f)+\beta_{1}(f)  \tag{3-1}\\
& =2 \pi \mathrm{fT} \ell+\beta_{1}(f) \tag{3-2}
\end{align*}
$$

where $\beta_{o}(f)$ is the phase shift due to the inherent transmission delay/unit length, $T$.
The delay/unit length is related to the velocity of propagation, $v$, and therefor, related to the inductance/unit length, $L$, and the capacitance/unit length, $C$.

$$
\begin{equation*}
\mathrm{v}=\frac{1}{\mathrm{~T}}=\frac{1}{\sqrt{\mathrm{LC}}}, \text { unit length/sec. } \tag{3-3}
\end{equation*}
$$

Also, the nominal or high frequency characteristic impedance, $R_{o}$, is dependent upon $L$ and $C$,

$$
\begin{equation*}
R_{o}=\sqrt{\frac{L}{C}} \text {, ohms. } \tag{3-4}
\end{equation*}
$$

Consequently, if $C$ is specified, then a knowledge of $v$ or $R_{o}$ determines $L$. Alternately, $v$ and $R_{o}$ may be specified which in turn determine $L$ and $C$,

$$
\begin{equation*}
L=\frac{R_{0}}{v} \tag{3-5}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathrm{C}=\frac{1}{\mathrm{VR}} . \tag{3-6}
\end{equation*}
$$

The relations (3-3) through (3-6) are useful in both the time domain and frequency domain modeling; furthermore, some manufacturers specify $v$ and $R_{o}$ rather than $C$ and $R_{0}$.

### 3.2 Measurement of the per Unit Length Resistance <br> and Capacitance, $R$ and $C$

Normally, $R$ and $C$ are specified by the manufacturer. If for some reason they must be determined experimentally, bridge measurements may be used. $R$ is determined by measuring the loop resistance, $\ell R$ using a D.C. resistance bridge, figs. 3.1 and 3.2. Similarly, $C$ is determined using a 1 kHz capacitance bridge. With a bridge having one terminal grounded, the measurement is straight-forward for a coaxial conductor unbalanced cable, fig. 3.3. However, for a shielded paired-conductor cable, in general, three capacitance measurements have to be made to determine the total cable capacitance, ${ }^{\ell} C_{A B}$,

$$
\begin{equation*}
\ell C_{A B}=\ell\left[C_{1}+\frac{C_{2} C_{3}}{C_{2}+C_{3}}\right] . \tag{3-7}
\end{equation*}
$$

The terminal designations are defined in fig. 3.4 while the equivalent circuit capacitances are defined in fig. 3.5. The measurement technique follows the procedure in [4] and consists of the following capacitance measurements: (1) $C^{\prime}$, A to ground with B grounded, (2) $C^{\prime \prime}$, A to ground with A and B shorted rogether, (3) $C^{\prime \prime \prime}, B$ to ground with $A$ grounded. The values of $\ell C_{1}, \ell C_{2}, \ell C_{3}$ are then given by

$$
\begin{align*}
& \ell C_{1}=\frac{C^{\prime}-C^{\prime \prime}+c^{\prime \prime \prime}}{2}  \tag{3-8}\\
& \ell C_{2}=\frac{C^{\prime}+C^{\prime \prime}-C^{\prime \prime \prime}}{2}  \tag{3-9}\\
& \ell C_{3}=\frac{-C^{\prime}+C^{\prime \prime}+C^{\prime \prime \prime}}{2} \tag{3-10}
\end{align*}
$$

### 3.3 Measurement of the Nominal Characteristic <br> Impedance, $\mathrm{R}_{\mathrm{o}}$

The nominal characteristic impedance $R_{o}$ is actually the limit of the high frequency impedance (2-33, 2-37) and is conceptually viewed as the characteristic impedance which depends solely upon the inductance and capacitance per unit length. Furthermore, $R_{o}$ is resistive, i.e., independent of $s$.

For large values $|s|$, the characteristic impedance is

$$
\underset{\mathrm{HF}}{\mathrm{Z}_{\mathrm{o}}(\mathrm{~s})}=\sqrt{\frac{\mathrm{L}}{\mathrm{C}}}[1+\delta(\mathrm{s})], \quad|\delta(\mathrm{s})| \ll \mid
$$

or

$$
\begin{equation*}
=R_{o}[1+\delta(s)] \tag{3-11}
\end{equation*}
$$

Now consider for the moment, that a step of current $I_{g} / s$ is applied to the impedance (3-11). The voltage across $Z_{o}(s)$ at $t=0^{+}$would jump up to $I_{g} R_{o}{ }_{0}$ volts and then commence to vary with time as dictated by the inverse Laplace transform of $I_{g} R_{o} \delta(s) / s$, fig. 3.6. By applying a step of voltage from a resistive source $R_{g}$ to the input of a long length of cable, it is possible to calculate the value of $R_{o}$ from the initial response of the sending end voltage $e_{s}(t)$, fig. 3.7.

First of all, when a signal is abruptly applied to a transmission line, the transmission line appears as an impedance equal to the characteristic impedance $Z_{o}(s)$. This is why the term surge impedance is also used in place of characteristic impedance. The impedance continues to appear as $Z_{o}(s)$ until reflections from the receiving end arrive back at the sending end. If the line is terminated in $Z_{0}(s)$ then there will be no reflections, and the line will continue to appear as an impedance equal to $Z_{o}(s)$. More will be said about such phenomena in the chapter on time domain reflectometry (Chapter 7).

Furthermore, for the initial response, only the high frequency components of $Z_{o}(s)$ contribute to the response. Consequently, for some time interval, say $t_{1}$, the high frequency approximation to $Z_{o}(s)$ (3-11) would be valid for predicting the initial response of the transmission line, as shown by the solid-line portion of $e_{s}(t)$ between 0 and $t_{1}$ in fig. 3.7. Also, in fig. $3.7 \mathrm{t}_{2}$ denotes qualitatively the time required for reflections to return to the sending end from the open circuited receiving end. The value of $t_{2}$ is set by $\ell$ and must be large enough so that reflections do not interfere with the observation of the abrupt rise in the neighborhood of $t=0^{+}$.

Consequently, the initial response of the cable to a voltage step of magnitude $E_{g}$ is given by

$$
\begin{align*}
E_{s}(s) & =\frac{E_{g}}{s} \frac{\left.z_{o}(s)\right\rfloor_{H F}}{\left.R_{g}+Z_{o}(s)\right]_{H F}}  \tag{3-12}\\
& =\frac{E_{g}}{s} \frac{R_{o}+R_{o} \delta(s)}{R_{g}+R_{o}+R_{o} \delta(s)} \tag{3-13}
\end{align*}
$$

or approximated by

$$
\begin{equation*}
E_{s}(s) \stackrel{E_{g}}{s} \frac{R_{o}}{R_{g}+R_{o}}[1+\delta(s)] \tag{3-14}
\end{equation*}
$$

because $R_{0} \gg R_{0}|\delta(s)|$ and $R_{0}+R_{g}$ is larger than $R_{0}$. In the time domain, the initial response of $e_{s}(t)$ consists of an abrupt jump $E$,

$$
\begin{equation*}
E=\frac{E_{g} R_{o}}{R_{g}+R_{o}} \tag{3-15}
\end{equation*}
$$

and then a slower upward trailing, fig. 3.7. From (3-15) $R_{o}$ is obtained as

$$
\begin{equation*}
R_{o}=\frac{E}{E_{g}-E} R_{g} \tag{3-16}
\end{equation*}
$$

If $R_{g}$ happens to equal $R_{o}$, then $E$ would equal $E_{g} / 2$. An example of an initial response expressible mathematically in closed form is given for $m=0.5$ in [5].

The measurement of $R_{o}$ can be implemented by the methods shown in figs. 3.8 and 3.9 for the coaxialconductor cable and the shielded paired-conductor cable, respectively. For the coaxial-conductor cable measurement, the effective generator impedance is $50 \Omega$.

For the paired-conductor case the generator impedance is $100 \Omega$ and the sampling oscilloscope display is set for channel A minus channel B which yields a display voltage

$$
\begin{equation*}
e_{d}(t)=e_{A}(t)-e_{B}(t)=e_{S}(t) \tag{3-17}
\end{equation*}
$$

Also, remember that the ground connection must be connected to the signal-ground, i.e., the zero potential for the balanced-transmission line (refer to fig. 3.5). For both cases, the cable length \& should be long enough so that reflections do not occur in the region of the initial response, fig. 3.7.

Closely related to the measurement of $R_{o}$ is the use of time domain reflectometry (TDR) to evaluate or test cables. The reader is referred to Chapter 7 of this report and in particular, Section 7.3 on a practical measurement method for observing the sending-end voltage.
3.4 Insertion Measurements for the Determination of the High Frequency Loss Parameter and the Exponent m

The loss parameter $K$ and the exponent $m$ are determined by model-fitting to data obtained from a time domain (pulse) insertion measurement, the basic method being shown in fig. 3.10. $e_{d 1}(t)$ and $e_{d 2}(t)$ are the observable quantities displayed on an oscilloscope before and after the insertion of the cable of length $\ell$ between a generator and matched load. The generator and load impedances are both equal to the nominal transmission line characteristic impedance. Denoting the inverse Laplace transform as $\mathcal{L}^{-1}\{ \}, e_{d 1}(t)$ and $e_{d 2}(t)$ are given by

$$
\begin{align*}
& e_{d 1}(t)=\mathcal{L}^{-1}\left\{T(s) \frac{E_{g}(s)}{2}\right\}  \tag{3-18}\\
& e_{d 2}(t) \quad \mathcal{L}^{-1}\left\{T(s) 2 E_{g}(s)\left[\frac{Z_{0}(s)}{R_{0}+Z_{o}(s)}\right]\left[\frac{R_{0}}{R_{0}+Z_{o}(s)}\right]\left[\frac{e^{-\ell \gamma(s)}}{1-[0(s)]^{2} e^{-2 \ell \gamma(s)}}\right]\right\} \tag{3-19}
\end{align*}
$$

where $T(s)$ is the oscilloscope transfer function (input to display), and $\tau(t)$ its inverse transform or impulse response. Also,

$$
\begin{equation*}
\rho(s)=\frac{R_{0}-Z_{0}(s)}{R_{0}+Z_{0}(s)} \tag{3-20}
\end{equation*}
$$

Aside from the factor $T(s)$ (3-19) is taken from (2-39), while (3-20) is from (2-41). The ratio, $E_{d 2}(s) / E_{d 1}(s)$, is equal to $E_{2}(s) / E_{1}(s)$,

$$
\begin{equation*}
\frac{E_{2}(s)}{E_{1}(s)}=4\left[\frac{Z_{0}(s)}{R_{0}+Z_{0}(s)}\right]\left[\frac{R_{0}}{R_{0}+Z_{o}(s)}\right]\left[\frac{e^{-\ell \gamma(s)}}{1-[\rho(s)]^{2} e^{-2 \ell \gamma(s)}}\right] \tag{3-21}
\end{equation*}
$$

because $T(s)$ cancels out along with the pulse dependence, $E_{g}(s)$. The ratio $E_{2}(s) / E_{1}(s)$ is equal to the scattering network parameter $\mathrm{S}_{21}(\mathrm{~s})$.

Briefly, the transmission line can be represented by the scattering parameters $\mathrm{S}_{11}(\mathrm{~s}), \mathrm{S}_{22}$ (s), and $S_{12}(s)=S_{21}(s)$. Figure 3.11 shows a transmission line inserted into another uniform transmission system of characteristic impedance $Z_{o}(s)$. In general, the incident and reflected voltages on the $Z_{o}(s)$ transmission line system are

$$
\begin{equation*}
v_{r 1}(s)=v_{11}(s) s_{11}(s)+V_{12}(s) s_{12}(s) \tag{3-22}
\end{equation*}
$$

and

$$
\begin{equation*}
v_{r 2}(s)=v_{11}(s) s_{21}(s)+v_{12}(s) s_{22}(s) \tag{3-23}
\end{equation*}
$$

The terminal voltages of the inserted transmission line are

$$
\begin{equation*}
\mathrm{v}_{1}(\mathrm{~s})=\mathrm{v}_{\mathrm{i} 1}(\mathrm{~s})+\mathrm{v}_{\mathrm{r} 1}(\mathrm{~s}) \tag{3-24}
\end{equation*}
$$

and

$$
\begin{equation*}
v_{2}(s)=v_{12}(s)+v_{r 2}(s) \tag{3-25}
\end{equation*}
$$

When the uniform transmission line is terminated in $Z_{0}(s)$, fig. $3.12, V_{12}(s)$ is zero, and the voltages become

$$
\begin{align*}
& v_{r 1}(s)=v_{i 1}(s) s_{11}(s)  \tag{3-26}\\
& v_{r 2}(s)=v_{11}(s) s_{21}(s)  \tag{3-27}\\
& v_{1}(s)=v_{i 1}(s)+v_{r 1}(s)  \tag{3-28}\\
& v_{2}(s)=v_{r 2}(s) \tag{3-29}
\end{align*}
$$

The voltage $V_{2}(s)$ corresponds to $E_{2}(s)$ in (3-21). If the inserted transmission line in fig. 3.12 is removed from the circuit and terminals 1 and 2 connected together, then $V_{r 1}(s)$ is zero, and

$$
\begin{equation*}
v_{1}(s)=v_{11}(s) \tag{3-30}
\end{equation*}
$$

which corresponds to $E_{1}(s)$ in (3-21). Then from (3-27), (3-29), and (3-30), there results

$$
\begin{equation*}
\frac{E_{2}(s)}{E_{1}(s)}=s_{21}(s) \tag{3-31}
\end{equation*}
$$

Returning now to consideration of (3-21), in terms of the model for high frequencies $\gamma(s)$ and $Z_{o}$ (s) become

$$
\begin{equation*}
\gamma(\mathrm{s})]_{\mathrm{HF}}=\mathrm{sT}[1+\delta(\mathrm{s})] \quad|\delta(\mathrm{s})| \ll 1 \tag{2-32}
\end{equation*}
$$

and

$$
\begin{equation*}
\left.Z_{o}(s)\right]_{H F}=R_{o}[1+\delta(s)] \quad|\delta(s)| \ll 1 \tag{2-33}
\end{equation*}
$$

which reduces (3-21) to

$$
\begin{equation*}
\left.\frac{E_{2}(s)}{E_{1}(s)}\right]=e^{-s T \ell-T \ell \delta(s)} \tag{3-32}
\end{equation*}
$$

or in terms of the Fourier transform

$$
\begin{equation*}
\left.\frac{E_{2}(j \omega)}{E_{1}(j \omega)}\right]_{H F}=e^{-j \omega T \ell-T \ell \delta(j \omega) .} \tag{3-33}
\end{equation*}
$$

Thus, the model attenuation at high frequencies

$$
\begin{equation*}
\left.\left.\alpha_{1}(\omega)\right]_{d B}=20 \log \left\lvert\, \frac{E_{1}(j \omega)}{E_{2}(j \omega)}\right.\right]\left._{H F}\left|=8.68 \frac{\ell K}{2 R_{o}}\right| \omega\right|^{m} \cos m \pi / 2 \tag{3-34}
\end{equation*}
$$

can be compared with the Fourier transform attenuation data computed from the experimentally observed waveforms, $e_{1}(t)$ and $e_{2}(t)$, in order to determine $K$ and $m$. As discussed in section 2.3 and illustrated in figs. $2.6-2.8$, this is accomplished in a tractable way using a $\log -\log$ graph, $\log \alpha_{1}(\omega)$ vs. $10 g$ $\omega$.

In practice, the insertion measurements are carried out using a 50 ohm sampling oscilloscope system which has two channels. Because $R_{0}$ may not be equal to 50 ohms or 100 ohms (for coaxial-conductor or paired-conductor cable, respectively) resistive impedance matching networks may be required, figs. 3.13 and 3.14. The matching networks are in an " L " configuration, each different value of $\mathrm{R}_{\mathrm{o}}$ requiring its own matching network. The networks used in this study are summarized in fig. 3.15.


Figure 3.1 Measurement of the total coaxial-conductor cable resistance, $\ell \mathrm{R}$.


Figure 3.2 Measurement of the total shielded paired-conductor cable resistance, $\ell$ R.


Figure 3.3 Measurement of the total coaxial-conductor cable capacitance, lC.


Figure 3.4 Terminal designations for measuring the total shielded pairedconductor cable capacitance, $\ell C$.


Figure 3.5 Shielded paired-conductor configurations and their equivalent capacitance circuits.

$\mathcal{L}\left[i_{g}(t)\right]=\frac{I_{g}}{s}$



Figure 3.6 The response of $\left.Z_{o}(s)\right]_{H F}$ to a current step of magnitude $I_{g}$.


Figure 3.7 The initial response (solid line) of the sending end voltage $e_{s}(t)$ due to a voltage step from a generator having a resistive source impedance.


Figure $3.8 \mathrm{R}_{\mathrm{o}}$ measurement set-up for a coaxial-conductor cable (unbalanced lines, generally).


Figure $3.9 \mathrm{R}_{\mathrm{o}}$ measurement set-up for a shielded paired-conductor cable (balanced lines, generally).


Figure 3.10 The time domain insertion method; $e_{d 1}(t)$ and $e_{d 2}(t)$ are the observables corresponding to $e_{1}(t)$ and $e_{2}(t)$, respectively, before and after the insertion of the cable of length $\ell . \tau(t)$ is the impulse response of the oscilloscope.


Figure 3.11 The scattering parameter representation of a transmission line inserted into a uniform transmission line system of characteristic impedance $Z_{o}(s)$.


Figure 3.12 Uniform transmission line system terminated in its characteristic impedance $Z_{o}(s)$.


Insertion plane

Figure 3.13 Practical implementation of the time domain insertion method for a coaxial-conductor cable.


Insertion plane

Figure 3.14 Practical implementation of the time domain insertion method for a shielded-paired-conductor cable.

(a)

(b)

| $\mathrm{R}_{\mathrm{o}}$ | $\mathrm{R}_{1}$ | $\mathrm{R}_{2}$ | $\mathrm{R}_{3}$ |
| :---: | :---: | :---: | :---: |
| 75 | 25 | 75 | 0 |
| 78 | 23.45 | 83.1 | 0 |
| 95 | 11.18 | 212.25 | 0 |
| 98 | 7.07 | 692.00 | 0 |
| 124 | 0 | 113.6 | 27.27 |
|  |  |  |  |

Figure 3.15 The resistive impedance matching networks, 50 ohms to $R_{o}$.
(a) Shielded paired-conductor cables,
(b) Coaxial-conductor cables.

## 4. AN OVERVIEW OF THE COMPUTER CALCULATION FOR THE TRANSMISSION LINE IMPULSE AND STEP RESPONSES

In this chapter a qualitative overview is presented which describes the computer calculations for the extraction of model parameters from measured data and the simulation of the impulse and step responses. The calculations are accomplished by the NBS Automatic Pulse Measurement System (APMS) [6,7].

### 4.1 Data Aćquisition Method

The data is generated by repetitive pulse excitation and is displayed by a sampling oscilloscope interfaced to and controlled by a minicomputer system, i.e., the NBS Automatic Pulse Measurement System (APMS). For a given display waveform, say $e_{d 2}(t)$, the data is acquired by making 500 observations at a fixed time position $t_{1}$ of 500 repetitions of the signal; the mean value [8,9] of the observations $\overline{e_{d 2}\left(t_{1}\right)}$ is calculated by the computer,

$$
\begin{equation*}
\overline{e_{d 2}\left(t_{1}\right)}=\frac{\sum_{r}^{500}\left[e_{d 2}\left(t_{1}\right)\right]_{r}}{500} \tag{4-1}
\end{equation*}
$$

The time window, $T_{W}$, is represented by 1024 uniformly spaced points; consequently, a set of 1024 mean values represent the waveform, fig. 4.1,

$$
\begin{equation*}
\hat{e}_{d 2}(n T)=\sum_{n=0}^{1023}\left\{\delta(t-n T) \overline{e_{d 2}(n T)}\right\} \tag{4-2}
\end{equation*}
$$

where $\delta(t)$ is the unit delta or impulse function multiplied by the strength $\overline{e_{d 2}(n T)}$ and where

$$
\begin{equation*}
\mathrm{T}=\mathrm{T}_{\mathrm{W}} / 1024 . \tag{4-3}
\end{equation*}
$$

Refer to figs. 4.5 and 4.6 for a graphical explanation of the delta function train (4.2). The waveforms $e_{d 1}(t)$ and $e_{d 2}(t)$ of the time domain insertion method, fig. 3.10, are both acquired by the procedure just described and will be denoted by $\hat{e}_{d 1}(n T)$ and $\hat{e}_{d 2}(n T)$, corresponding to the form of (4.2).

$$
\text { 4.2 Fourier Transformation of } \hat{e}_{\mathrm{d} 1}(\mathrm{nT}) \text { and } \hat{\mathrm{e}}_{\mathrm{d} 2}(\mathrm{nT})
$$

The discrete acquired waveforms $\hat{e}_{d 1}(n T)$ and $\hat{e}_{d 2}(n t)$ are each transformed to the frequency domain using the Fast Fourier Transform (FFT) algorithm for the discrete Fourier transformation [10]. To do so, for each waveform a 2048 point time window is used which includes 1024 points of zero value, fig. 4.2. The extension of the waveform to 2048 points increases the frequency resolution by decreasing the spacing between the discrete frequencies to one half of that available from the unaugmented waveform, fig. 4.3 and 4.4. The frequency domain spacing $\Delta f$ is equal to $1 / T_{W}$ and the frequency window, $\mathrm{F}_{\mathrm{W}}$ remains unchanged when both N and $\mathrm{T}_{\mathrm{W}}$ are doubled, but the number of points in $\mathrm{F}_{\mathrm{W}}$ are doubled.

The result of the FFT operation yields the discrete functions $\hat{\mathrm{E}}_{\mathrm{d} 1}\left(\frac{\mathrm{k}}{2 \mathrm{NT}}\right)$ and $\hat{\mathrm{E}}_{\mathrm{d} 2}\left(\frac{\mathrm{k}}{2 \mathrm{NT}}\right)$ which represent complex values at the discrete frequencies $k / 2 N T$, e.g., the discrete waveform, $\hat{e}_{d 2}$ (nT), of fig. 4.2 transforms to

$$
\begin{equation*}
\hat{E}_{d 2}\left(\frac{k}{2 N T}\right)=\sum_{n=0}^{2 N-1} \hat{e}_{d 2}(n T) e^{-j 2 \pi n k / 2 N} ; k=0,1, \ldots 2 N-1 \tag{4-4}
\end{equation*}
$$

where $N$ equals 1024; consequently, $\hat{E}_{d 2}(k / 2 N T)$ consists of 2048 complex values, one value for each of the 2048 discrete frequencies, $k / 2 N T$. Furthermore, $\hat{E}_{d 2}(k / 2 N T)$ is a symmetrical periodic function and as such contains on1y 1024 distinct values.

$$
\text { 4.3 The Calculation of } \hat{S}_{21}(k / 2 N T) \text { and } \hat{\alpha}_{1}(k / 2 N T)
$$

The discrete insertion ratio $\hat{\mathrm{S}}_{21}(\mathrm{k}, 2 \mathrm{NT})$ is computed in the same manner as its continuous counterpart,

$$
\begin{equation*}
S_{21}(j \omega)=\frac{E_{d 2}(j \omega)}{E_{d 1}(j \omega)}, \tag{3-31}
\end{equation*}
$$

hence

$$
\begin{align*}
\hat{\mathrm{S}}_{21}(k / 2 N T) & =\frac{\hat{\mathrm{E}}_{\mathrm{d} 2}(k / 2 N T)}{\hat{\mathrm{E}}_{\mathrm{d} 1}(k / 2 N T)}  \tag{4-5}\\
& =\left|\hat{S}_{21}(k / 2 N T)\right| / \hat{\beta}_{1}(k / 2 N T) \tag{4-6}
\end{align*}
$$

where $k=0,1, \ldots 2 N-1$, and $\left|\hat{S}_{21}\right|$ and $\hat{\beta}_{1}$ are the magnitude and phase of the discrete function $\hat{S}_{21}$ at the discrete frequencies $\mathrm{k} / 2 \mathrm{NT}$, respectively.

However, an averaging process is used in arriving at $\hat{S}_{21}(k / 2 N T)$. The process is repeated using six different acquisitions for $\hat{E}_{d 2}(k / 2 N T)$ and a single acquisition for $\hat{E}_{d 1}(k / 2 N T)$. Then, the resultant $\hat{\mathrm{S}}_{21}(\mathrm{k} / 2 \mathrm{NT})$ data sets are expressed in terms of the discrete attenuation $\hat{\alpha}_{1}(k / 2 N T)$ and phase $\hat{\beta}_{1}(k / 2 N T)$,

$$
\text { Attenuation, } \hat{\alpha}_{1}(k / 2 N T) \equiv 20 \log \left|\hat{\mathrm{~S}}_{21}(\mathrm{k} / 2 \mathrm{NT})\right|^{-1}, \mathrm{~dB}
$$

$$
\text { Phase, } \hat{B}_{1}(k / 2 N T)=-\operatorname{Tan}^{-1}\left\{\frac{\text { Imag. Pt. } \hat{S}_{21}^{(k / 2 N T)}}{\text { Real Pt. } \hat{S}_{21}^{(k / 2 N T)}}\right\} \text {, degrees. }
$$

and the mean values are determined by

$$
\begin{align*}
& \overline{\hat{\alpha}_{1}(k / 2 N T)}=\frac{\sum_{r}^{6} 20 \log \left|\hat{S}_{21}(k / 2 N T)\right|_{r}^{-1}}{6}, d B  \tag{4-7}\\
& \overline{\hat{\beta}_{1}(k / 2 N T)}=\frac{\sum_{r}^{6}\left[\hat{\beta}_{1}(k / 2 N T)\right]}{6}, \text { degrees } \tag{4-8}
\end{align*}
$$

where

$$
\begin{equation*}
\left[\hat{\mathrm{S}}_{21}(\mathrm{k} / 2 \mathrm{NT}]_{\mathrm{r}}\right]^{-1}=\frac{\left[\hat{\mathrm{E}}_{\mathrm{d} 1}(\mathrm{k} / 2 \mathrm{NT})\right]}{\left[\hat{\mathrm{E}}_{\mathrm{d} 2}(\mathrm{k} / 2 \mathrm{NT})\right]_{\mathrm{r}}} \tag{4-9}
\end{equation*}
$$

$$
\begin{equation*}
=\left|\hat{\mathrm{S}}_{21}(\mathrm{k} / 2 \mathrm{NT})\right|_{\mathrm{r}}^{-1}<-\left[\hat{\beta}_{1}(\mathrm{k} / 2 \mathrm{NT})\right]_{\mathrm{r}} . \tag{4-10}
\end{equation*}
$$

and $\left[\hat{E}_{d 2}(k / 2 N T)\right]_{r}$ corresponds to the $r-t h$ acquisition from $e_{d 2}(t)$, i.e., $\left[\hat{e}_{d 2}(n T)\right]_{r}$.
The discrete values $\overline{\hat{\alpha}_{1}(k / 2 N T)}$ and $\overline{\hat{\beta}_{1}(k / 2 N T)},(4-7)$ and (4-8), are computed for the first three hundred harmonics, because in most cases 300 harmonics cover a magnitude range of about 0 to -30 dB . The capability for any other number of harmonics is retained in the computer program. At each value of frequency, the standard deviation [8] for (4-7) and (4-8) are computed as a measure of the consistency of the data. In particular, the standard deviation of the attenuation, $\sigma_{s}$, is given by

$$
\begin{equation*}
\sigma_{s}(k / 2 N T)=\left[\frac{\sum_{r}^{6}\left[\overline{\hat{\alpha}_{1}(k / 2 N T)}-\hat{\alpha}_{1}(k / 2 N T)_{r}\right]^{2}}{5}\right]^{1 / 2} . \tag{4-11}
\end{equation*}
$$

Past experience has shown that consistent data yields $\sigma_{s}$ values in a typical range of 0.02 to 0.2 dB for an attenuation range of one to 30 dB .

### 4.4 The Extraction of the Model Parameters $K$ and m

The parameters $K$ and $m$ are obtained by fitting the model at high frequencies to the (mean) attenuation data $\left[\overline{\hat{\alpha}_{1}(k / 2 N T)}\right]$ as was pointed out in the discussion of (3-34) and figs. 2.6-2.8. The values of $K$ and $m$ of the equation

$$
\begin{equation*}
\alpha_{1}(\omega)=\left(\frac{8.68 \ell \frac{\cos m \pi / 2}{2} R_{o}}{R_{0}|\omega|^{m} .{ }^{m} .}\right. \tag{3-34}
\end{equation*}
$$

can be determined from the straight-1ine parameters of the equation resulting from the $\log$ of (3-34),

$$
\begin{equation*}
\log \quad \alpha_{1}(\omega)=\log \left(\frac{8.68 \ell \cos m \pi / 2}{2 R_{0}}\right)+\log K+m \log |\omega| ; \tag{4-12}
\end{equation*}
$$

that is,

$$
\begin{equation*}
\log \hat{\alpha}_{1}(\mathrm{k} / 2 \mathrm{NT})=\log \left(\frac{8.68 \ell \cos m \pi / 2}{2 R_{\mathrm{o}}}\right)+\log K+m \log |\mathrm{k} / 2 \mathrm{NT}| \tag{4-13}
\end{equation*}
$$

The parameters $\log \mathrm{K}$ and m can be determined from the straight-1ine parameters of the logarithm of the mean attenuation $(4-7), \log \left[\overline{\hat{\alpha}_{1}(k / 2 N T)}\right]$.

The slope $m$ is determined by calculating the log-log graph slope from two values of the graph separated by 50 abscissa points, e.g., $k=50$, and $k=100$. The slope so determined is designated $m_{1}$. Next $K$ is calculated using $m$, in (4-13) from

$$
\begin{equation*}
\log \hat{\alpha}_{1}(k / 2 N T)=\log \left(\frac{8.68 \ell \cos m_{1} \pi / 2}{2 R_{o}}\right)+\log K+m_{1} \log |k / 2 N T| \tag{4-14}
\end{equation*}
$$

by an iterative solution method; the result is designated $K_{1}$.
The process for calculating $m_{1}$ and $K_{1}$ is repeated using the next two points $k=51$ and $k=101$; the results are designated $m_{2}$ and $K_{2}$. The process is repeated 150 times ending with $k=200$ and $k=250$. The mean values of $m_{r}$ and $K_{r}$ are then determined as

$$
\begin{equation*}
\overline{\mathrm{m}}=\sum_{\mathrm{r}=1}^{150} \mathrm{~m}_{\mathrm{r}} / 150 \tag{4-15}
\end{equation*}
$$

and

$$
\begin{equation*}
\bar{K}=\sum_{r-1}^{150} K_{r} / 150 \tag{4-16}
\end{equation*}
$$

The calculation process is started at $k=50$, rather than $k=0$ to avoid the small errors which appear at the beginning of the frequency window due to system errors, noise, etc.

Finally, the entire process is repeated three times; starting from $e_{d 1}(t)$ and $e_{d 2}(t)$ the complete operations of sections $4.1,4.2$, and 4.3 are repeated to yield three independent functions $\left[\widehat{\alpha}_{1}(k / 2 N T)\right] 1,\left[\bar{\alpha}_{1}(k / 2 N T)\right]_{2}$, and $\left[\hat{\alpha}_{1}(k / 2 N T)\right]_{3}$ from which three sets of mean values for $m$ and $k$ are determined:

$$
\begin{aligned}
& {\left[\overline{\hat{\alpha}_{1}(\mathrm{k} / 2 \mathrm{NT})}\right]_{1} \longrightarrow[\overline{\mathrm{~m}}]_{1}, \quad\left[\overline{\mathrm{~K}}_{1}\right]_{1}} \\
& {\left[\hat{\alpha}_{1}(\mathrm{k} / 2 \mathrm{NT})\right]_{2} \longrightarrow[\overline{\mathrm{~m}}]_{2}, \quad[\overline{\mathrm{~K}}]_{2}} \\
& {\left[\hat{\alpha}_{1}(\mathrm{k} / 2 \mathrm{NT})\right]_{3} \longrightarrow[\overline{\mathrm{~m}}]_{3}, \quad[\overline{\mathrm{~K}}]_{3} .}
\end{aligned}
$$

Then, the mean values for $[\overline{\bar{W}}]$ and $[\bar{K}]$ are calculated, where

$$
\begin{align*}
& {[\overline{\bar{m}}]=\frac{[\overline{\mathrm{m}}]_{1}+[\overline{\mathrm{n}}]_{2}+[\overline{\mathrm{m}}]_{3}}{3}}  \tag{4-17}\\
& \overline{\overline{\mathrm{~K}}]}=\frac{[\overline{\mathrm{K}}]_{1}+[\overline{\mathrm{K}}]_{2}+[\overline{\mathrm{K}}]_{3}}{3} \tag{4-18}
\end{align*}
$$

and the resultant values are the ones used for the model parameters $m$ and $K$. The standard deviations, $\sigma_{\mathrm{m}}$ and $\sigma_{\mathrm{K}}$ are also calculated for control purposes to insure that the data are consistent. They are

$$
\begin{equation*}
\left.\sigma_{m}=\left[\sum_{\mathrm{r}=1}^{3} \frac{\{\overline{[\overline{\mathrm{~m}}]}-[\overline{\mathrm{m}}]}{2}\right\}^{2}\right]^{1 / 2} \tag{4-19}
\end{equation*}
$$

and

$$
\begin{equation*}
\sigma_{K}=\left[\sum_{\mathrm{r}=1}^{3} \frac{\left\{\overline{\bar{K}]}-[\overline{\mathrm{K}}]_{\mathrm{r}}\right\}^{2}}{2}\right]^{1 / 2} . \tag{4-20}
\end{equation*}
$$

The impulse response of a cable of length $\ell$ doubly terminated in its nominal characteristic impedance $R_{o}$ is obtained from (2-39) with $E_{g}(s)=1$,

$$
\begin{equation*}
H(s)=2\left[\frac{Z_{0}(s)}{R_{0}+Z_{o}(s)} \cdot\right]\left[\frac{R_{0}}{R_{0}+Z_{o}(s)}\right]\left[\frac{e^{-\ell \gamma(s)}}{1-[\rho(s)]^{2} e^{-2 \ell \gamma(s)}}\right] \tag{4-21}
\end{equation*}
$$

where

$$
\begin{align*}
Z_{o} & =\left[\left(R+s L+K s^{m}\right) / s C\right]^{1 / 2} \\
\gamma(s) & =\left[\left(R+s L+K s^{m}\right) s c\right]^{1 / 2} \\
\rho(s) & =\frac{R_{o}-Z_{o}(s)}{R_{o}+Z_{o}(s)} \tag{2-41}
\end{align*}
$$

$$
\infty
$$

The received voltage $(4-21)$ due to a unit impulse for $e_{g}(t), E_{g}(s)=1$, is equal to one half of $E_{2}(s) / E_{1}(s)$, (3-21); consequently, from (3-31) the impulse response is

$$
\begin{equation*}
H(s)=\frac{1}{2} S_{21}(s) \tag{4-22}
\end{equation*}
$$

Using the model parameters ( $\mathrm{R}, \mathrm{L}, \mathrm{K}, \mathrm{m}$, and C C ) which are now known, a discrete 1024 point representation of the insertion function $\frac{1}{2} S_{21}(s)$ is constructed. Next, the inverse discrete Fourier transformation of $\frac{1}{2} \hat{S}_{21}(k / N T)$ will yield the discrete time-domain impulse response,

$$
\begin{equation*}
\hat{h}(n T)=\frac{1}{N} \sum_{k=0}^{N-1} \frac{1}{2} \hat{S}_{21}(k / N T) e^{j 2 \pi n k / N} ; \quad n=0,1, \ldots N-1 \tag{4-23}
\end{equation*}
$$

A typical display for $\hat{h}(n T)$ would be similar to $\hat{e}_{d 2}(n T)$, fig. 4.1.

### 4.6 Determination of the dc Level for the Impulse Response, $\hat{h}(n T)$

The dc level of the impulse response must be established before the step response can be calculated. Theoretically, the point $n=0$ should yield $\hat{h}(o)=0$; if such were the case, then by setting the level of $\hat{h}(n T)$ equal to zero at $n=0$, the dc level would be properly located. However, due to the system errors in the region of $o \leq n<10$, the zero level can not be determined. Similar effects are encountered at the end of the time window near $n=1024$.

Physically, the impulse response of the cable at the end of the time window is not equal to zero; it will be a small value, $10^{-4}$, (typically $10^{-3}$ of the peak value of the impulse response). Consequently, to set the dc level, the voltage at a specified time, say $n_{1}$, is calculated using the model and the identified model parameters. To do so, the Fourier components of the transfer function (impulse response) are summed for a specified time, typically, at $90 \%$ of the time window, $0.9 \mathrm{~T}_{\mathrm{w}}$. At the present time on the APMS this sumation of the Fourier series requires considerable computation time, e.g.,
several hours, because a large number of harmonics must be summed withgut recourse to array processing computation. Typically, for a 1 KHz separation of harmonics $5.12 \times 10^{5}$ harmonics are required for a one microsecond time window containing 1024 data points (harmonics from 1 kHz to 512 MHz ). The typical convergence at $0.9 \mathrm{~T}_{\mathrm{w}}$ is such that the voltage converges to about $10^{-3} \mathrm{~V}_{\max }$ with a peak oscillatory variation in the range of $10^{-4} \mathrm{v}_{\max }$ to $10^{-5} \mathrm{v}_{\max }$.

Once the voltage for $n_{1}$ has been computed, then the dc level of the impulse response waveform, $\hat{h}(n T)$, is shifted so that the value of $\hat{h}\left(n_{1} T\right)$ is equal to the computed voltage at $n_{1}$.

### 4.7 Simulation of the Step Response

The impulse response (4-23) may be expressed in the form

$$
\begin{equation*}
\hat{h}(n T)=\sum_{n=0}^{1023} h(n T) \delta(t-n T) \tag{4-24}
\end{equation*}
$$

where

$$
\begin{equation*}
T=T_{W} / 1024 \tag{4-25}
\end{equation*}
$$

and each term for the product $h(n T) \delta(t-n T)$ is given by a corresponding sum over $k$ in (4-23).
The discrete step response is obtained by integrating (4-24) and then sampling the result to provide discrete data values. Integrating the impulse train (4-24) gives

$$
\begin{equation*}
f(t)=\int \hat{h}(n T) d t=T \sum_{m=0}^{n} h(m T) u(t-m T) \tag{4-26}
\end{equation*}
$$

where $u(t)$ is the unit step function. (4-26) is a staircase waveform whose envelope is the integral of the impulse train. By sampling $f(t)$ at the points $t=n T$, the discrete waveform for the step response is obtained as

$$
\begin{align*}
\hat{f}(n T) & =T \sum_{n=0}^{1023} f(n T) \delta(t-n T)  \tag{4-27}\\
& =T \sum_{n=0}^{1023} \delta(t-n T) \quad \sum_{m=0}^{n} h(m T) . \tag{4-28}
\end{align*}
$$

The integration of the impulse response (4-24) to yield the step response (4-28) is illustrated graphically in figs. 4-5 through 4-8.


Figure 4.1 A given display waveform is represented by 1024 data points, each point being the mean value of 500 observations.


Figure 4.2 A given display waveform of 1024 points augmented by 1024 zeros to yield a total time window of 2048 points.


Figure 4.3 One cycle of the periodic 1024 discrete Fourier transform corresponding to $\hat{e}_{d 1}(n T)$, fig. 4.1.


Figure 4.4 One cycle of the pericdic 2048 point discrete Fourier transform corresponding to the augmented-zero waveform, fig. 4.2.


Figure 4.5 The sequence of values representing the discrete impulse response $\hat{h}(n T)$.


Figure 4.6 The sequence $\hat{h}(\mathrm{nT})$ represented as a train of delta functions of varying strengths. Keep in mind that each delta function is infinite in amplitude while its strength is given by its integral. This representation corresponds to (4.24) and is the mathematical representation for a sequence of values.


Figure 4.7 The integral of the delta function train, fig. 4.6; yields a staircase function. Each step along the staircase corresponds to the strength of the corresponding impulse in fig. 4.6. $f(t)$ is a piece-wise continuous function, (4-26).


Figure 4.8 The sequence of values representing the discrete step response $\hat{\mathrm{f}}(\mathrm{uT})$. This is the discrete function corresponding to the piece-wise continuous function (4-28) shown in Fig. 4.7.

## 5. TYPICAL RESULTS FOR TRANSMISSION CHARACTERIZATION AND RESPONSE

In this chapter typical results are presented for the determination of the transmission line parameters and the corresponding frequency and time domain responses. These results are typical of the data obtained for all of the cables considered in this report; data on specific cable types vs. length are given in Appendix A. Here, a single cable type of fixed length is used to illustrate results for parameter characterization and response simulation.

### 5.1 Description of the Cable

The cable results reported below were obtained using a 320 meter ( 1050 ft .) length of shielded paired-conductor transmission line. This cable is commercially available and for the purposes of this report is designated by the letter $I$. The nominal specifications for the cable are as follows:

```
Nominal characteristic impedance - 124 ohms
Overall diameter - 1.07 cm (0.420 in.)
Inductance/unit length - 620 nH/m (189 nH/ft)
Capacitance/unit length - 40.3 pf/m (12.3 pf/ft)
Resistance/unit length - 61.7 \times 10 -3 ohms/m (18 ohms/ft.)
```


### 5.2 The Results

The results are divided into three groups. The first one illustrates the acquired time domain insertion voltages and the subsequent computed (FFT) complex frequency domain data for the insertion ratio $S_{21}(j \omega)$. The second section shows the model simulated responses obtained. In the third section the simulation results for the impulse and step responses are shown.

Figures 5-1 and 5-2 show the acquired waveforms $\hat{e}_{d 1}(n T)$ and $\hat{e}_{d 2}(n T)$ which are the discrete insertion waveforms corresponding to equations (3-18) and (3-19); they were acquired in the manner discussed in Section 4.1. Next, the discrete Fourier transforms of $\hat{e}_{d 1}(n T)$ and $\hat{e}_{d 2}(n T)$ are computed and are used to calculate the complex ratio $\left[\hat{\mathrm{E}}_{\mathrm{d} 2}(\mathrm{k} / 2 \mathrm{NT})\right] /\left[\hat{\mathrm{E}}_{\mathrm{d} 1}(\mathrm{k} / 2 \mathrm{NT})\right]$ which is the insertion ratio $\hat{\mathrm{S}}_{21}(\mathrm{k} / 2 \mathrm{NT})$, Sections 4.2 and 4.3. Table 5-1 lists the complex data for $\hat{S}_{21}(k / 2 N T)$, (4.10), along with the standard deviations of each value of $\left|\hat{S}_{21}\right|$ and the phase $B_{1}$. Figure 5-3 shows a log-log graph of the attenuation $\alpha(\mathrm{k} / 2 \mathrm{NT})$, (4.7). The phase data in Table 5-1 is used only to ascertain that the data for $\hat{\mathrm{S}}_{21}(\mathrm{k} / 2 \mathrm{NT})$ possesses a small enough standard deviation in the phase data, which in conjunction with a small standard deviation for the magnitude $\left|\hat{S}_{21}\right|$, indicates that the data is consistent.

The parameters $m$ and $K$ are extracted from the data shown in fig. 5-3 according to the theory in Section 3.4 and by the method of Section 4.4. Using the extracted values of $m$ and $K$, the simulated cable attenuation $\hat{a}_{1}(k / 2 N T)$ and its associated minimum phase shift $\hat{\beta}_{1}(k / 2 N T)$ are calculated, figs. 5-4 and 5-5.

By taking the inverse discrete Fourier transform, using the inverse FFT algorithm, of the frequency domain simulation data, the simulated impulse response of the cable, h(nT), is obtained, fig. 5-6. The D. C. level of the impulse response has been carefully established by the procedure described in section 4.6. To validate the (simulated) impulse response, the acquired reference waveform $\hat{e}_{d 1}$ ( nT ) is convolved with the impulse response and the result is compared to the acquired insertion waveform $\hat{e}_{\mathrm{d} 2}$ ( nT ). Figure 5-7 shows $\hat{e}_{d 2}(\mathrm{nT})$, fig. $5-2$, with an expanded time scale while fig. $5-8$ shows the result of convolving $\hat{e}_{d 1}(n T)$, fig. 5-1 with the simulated impulse response $\hat{h}(n T)$, fig. 5-6. Comparison of figs. 5and 5-8 shows no significant differences in shape, hence the model parameters accurately characterize the cable.

The step response of the cable, $\hat{f}(n T)$, is obtained by integration of the impulse response $\hat{h}(n T)$, fig. 5-6, as discussed in Section 4.7; the result is shown in fig. 5-9. Note that step response $\hat{f}(\mathrm{nT}$ ) has reached the normalized value of about 0.86 at the end of the time window, 999 nanoseconds. According to eq. (2-52) the final value referred to the sending end generator voltage $E_{g}$ would be $R_{o} /\left(2 R_{o}+\ell R\right)$ or 0.463 for this particular cable (parameter values given in Section 5.1). Without losses at dc ( $R=0$ ), this value would be 0.5. Normalizing the lossless value to 1.0 , gives a final value dc value, i.e., the value for $t=\infty$ ) of 0.926 . Consequently, it is seen that the cable response has not reached the true final value of 0.926 in one microsecond. Consequently, the transmission line could not accurately transmit pulses changing state every microsecond.

Figure 5.1 The acquired waveform $\hat{e}_{d I}(\mathrm{nT})$ for 320.04 meters ( $1050 \mathrm{ft}$. ) of cable I .

Figure 5.2 The acquired waveform $\hat{e}_{d 2}(n T)$ for 320.04 meters ( 1050 ft. ) of cable $I$.

100.0 MHz
$\hat{e}_{d 2}(\mathrm{nT})$
Frequency $\longrightarrow$
1.0 MHz
10.0 MHz for 320.04 meters ( $1 u$ ùu $f$. .) of cable I.
1.0 MHz

10.0 MHz
100.0 MHz

Frequency $\longrightarrow$



150.0 mV
75.0 mV
0.0 mV
0.0
$\hat{\mathrm{h}}(\mathrm{nT})$, Fig. 5.6 for 320.04 meters ( 1050 ft .) of cable I.
Figure 5.8 The result of convolving $\hat{e}_{d 1}(\mathrm{nT})$, Fig. b.1, with the simulated impulse response,
150.0 m
蹅
500 ns
1000 ns
$1.0^{\circ} \mathrm{V}$
0.5 v
0.0
s7TOM

including the standard deviations of the magnitude and phase.

List of magnitude and phase for the insertion ratio $\hat{S}_{21}(\mathrm{k} / 2 \mathrm{~N})$
including the standard deviations of the magnitude and phase.
NUMEER OF NEASUREMENTS=


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## 6. USE OF TIME DOMAIN MEASUREMENTS TO ESTIMATE MAXIMUM CABLE BIT RATE AND BIT-ERROR-RATE DEGRADATION

This chapter is concerned with developing a method for using the time domain measurement techniques and data presented in the previous five chapters to obtain estimates of two parameters of great interest to the digital communication system designer. These are (1) the maximum bit rate that may be transmitted through a given length and type of cable, and (2) the bit-error-rate that may be expected for a given system signal-to-noise ratio. It will be shown that a rather unusual plot of the cable step response may be used to estimate these two parameters.

### 6.1 Digital Communications System Model

Since the purpose of this study is to characterize certain transmission lines (as opposed to complete communication systems) for digital communication applications, a simple idealized transmitter/ receiver model has been chosen. By avoiding more complex "realistic" system models it becomes simpler to observe the deleterious effects of the transmission line, alone, in an otherwise perfect system.

The chosen system model is shown in fig. 6.1. The following important assumptions have been made with respect to this transmitter/receiver model:
a. The transmitter generates a binary non-return-to-zero (NRZ) signal of perfect rectangular pulse shape. That is, a "zero" is represented by a constant zerovolts output, a "one" by a constant V-volts output and the transitions from a "zero" to a "one" or vice versa are instantaneous. The transmitter bit rate is controlled by the trigger generator.
b. The receiver is a perfect, triggerable threshold detector such that at any time $t$ it may be triggered to sample the input signal and decide if the input voltage is less than or greater than some voltage $V_{T H}$. If the receiver decision is "less than $\mathrm{V}_{\mathrm{TH}}$ ", then a "zero" is assumed to have been received; and if the decision is "greater than $\mathrm{V}_{\mathrm{TH}}$ ", then a "one" is assumed to have been received.
c. The trigger generator and variable delay are assumed to allow perfect timing control between the transmitter and receiver for any bit rate. Thus, compensation for any cable delay is provided and the delay may be "fine tuned" to allow a receiver sample to be taken at any chosen time within a bit period.
d. The transmitter output impedance, the receiver input impedance and the summing circuit's input and output impedance are matched to the high frequency limit of the cable characteristic impedance; i.e., $\sqrt{\mathrm{L} / \mathrm{C}}$ (see eqn. 2-35).
(The noise source will be discussed in Section 6.4 ; for the discussion in this section switch $S_{1}$ is assumed to be open.)

### 6.2 Maximum Bit Rate Estimation

From the model assumptions it is evident that the transmitter/receiver pair, connected directly together, represents a system with no bandwidth (or, equivalently, bit rate) limitation. That is, since the transmitter voltage transitions are assumed to be instantaneous, and the receiver has no speed of sampling restrictions, then the system can operate perfectly at any bit rate. In addition, it should be noted that this system is not code dependent; it will work perfectly for any particular sequence of "ones" and "zeros" generated by the transmitter.

If a section of transmission line is now connected into the system at the insertion plane indicated in fig. 6.1, two changes in the operation of the system occur; the system becomes bit rate limited and code dependent. In order to analyze and quantize these cable-induced system limitations, the example measurements on 1050 ft . of cable I presented in Chapter 5 will be extended

The step response for cable $I$ for $0 \leq t<1 \mu s$ is shown in fig. 6.2. This, of course, is the voltage that would be observed at the receiver if a "one" were transmitted preceded by an infinite (in theory) string of "zeros". At the end of the time window the voltage is only about 0.86 V and, as was shown in Chapter 5, will rise to only 0.926 V as $\mathrm{t} \rightarrow \infty$. The fact that the received voltage never reaches V volts is due to the DC resistive losses in the cable.

If the situation were reversed such that a "zero" were sent at $t=0$, preceded by an infinite string of "ones", the response voltage seen at the receiver would be as shown in fig. 6.3. That is, it would be 0.926 V at $\mathrm{t}=0$, decay to 0.066 V at $\mathrm{t}=1 \mu \mathrm{~s}(0.926 \mathrm{~V}-.0 .86 \mathrm{~V}$ ) and finally reach 0 V as $t \rightarrow \infty$.

From this information the receiver may now be adjusted for optional operation. That is, a threshold level, $\mathrm{V}_{\text {TH }}$, and a sampling time $\mathrm{t}_{\mathrm{s}}$ may be chosen to maximize the receiver's ability to accurately detect whether a "one" or a "zero" was transmitted. In both the "zero"-to-"one" and "one"-to-"zero" transitions mentioned above the received voltage varied between the limits of 0 V and 0.926 V in the interval $0 \leq t \leq \infty$. Thus a logical choice for $V_{T H}$ is $V_{f} / 2$ where $V_{f}$ is a function of the type and length of cable and, as was discussed in Chapter 5, is calculated as

$$
\begin{equation*}
v_{f}=\frac{2 R_{o}}{2 R_{o}+\ell R} \tag{6-1}
\end{equation*}
$$

For the cable I parameters listed in Chapter 5 ( $\mathrm{R}_{\mathrm{o}}=124 \Omega, \ell=1050 \mathrm{ft}$. and $\mathrm{R}=18.8 \mathrm{~m} \Omega / \mathrm{ft}$ ) substitution into eqn. $6-1$ yields $\mathrm{V}_{\mathrm{f}}=0.926 \mathrm{~V}$ and thus $\mathrm{V}_{\mathrm{TH}}=0.463 \mathrm{~V}$. Any other choice of $\mathrm{V}_{\mathrm{TH}}$ will bias the receiver such that it no longer would have an equal ability to detect a "one" and a "zero".

The logical choice for a receiver sampling time is at the end of a bit period (note that this choice applies only to this model, e.g., no time jitter). From fig. 6.2 it is observed that in the time window $0 \leq t \leq 1 \mu_{s}$ the "one" voltage is maximum at $t=1 \mu \mathrm{~s}$. Also, from fig. 6.3, the "zero" voltage is minimum at $t=1 \mu_{s}$. Therefore, since the received "one" or "zero" voltage is either rising or decaying monotonically, the best time to sample is always at the end of a bit period, regardless of tie actual value of the bit period.

Figure 6.4 is a superposition of the positive and negative step responses shown in figs. 6.2 and 6.3. As expected, the two curves intersect at $V_{T H}=V_{f} / 2$ which occurs at $t \approx 50 \mathrm{~ns}$. If the system were operated at the bit rate corresponding to this intersect time; i.e., $\mathrm{BR}=1 / 50 \mathrm{~ns}=20 \mathrm{Mb} / \mathrm{s}$, then the receiver voltage at the sample time would be $\mathrm{V}_{\mathrm{TH}}$ regardless of whether a "one" or a "zero" were sent. Furthermore, if the system were operated at any faster rate, say at $B R=1 / 10 \mathrm{~ns}=100 \mathrm{MB} / \mathrm{s}$, then the receiver would always be wrong because a "one" signal (or a "zero" signal) would not have had sufficient time to reach a value above (or below) $\mathrm{V}_{\mathrm{TH}}$ • Also, if the system were operated at any slower rate, say at $B R=1 / 100 \mathrm{~ns}=10 \mathrm{Mb}$, the receiver would always be correct because a "one" signal (or a "zero" signal) would have had sufficient time to reach a value above (or below) $\mathrm{V}_{\mathrm{TH}}$.

Thus it is observed that the insertion of the cable into the otherwise unlimited-bit-rate system imposes a maximum bit rate limitation, and also that the value of this maximum bit rate may be obtained
as the reciprocal of the intersect time of the positive and negative cable step responses.
It is admittedly difficult to extract receiver voltage versus bit rate information from fig. 6.4. But by redrawing the positive and negative cable step response voltages on an $X$-axis scaling of $\log _{10}$ (bit rate) instead of time it becomes a simple matter to relate the intersect point to a maximum bit rate. This plot is shown in fig. 6.5 and represents the desired objective of this section. From it the maximum permissible bit rate for correct system operation may be read directly by observing the point on the X -axis, or bit rate axis, at which the two curves intersect, in this case about $20 \mathrm{Mb} / \mathrm{s}$.

Since this plot (fig. 6.5) is central to the remaining developments of this chapter, a few additional remarks about it are in order. The data from which it is derived, the positive and negative cable step responses of fig. 6.4, consist of two discrete time signals, each of 1024 points in a 1 رs window. That is, the step responses are really a set of 1024 voltage points spaced $1 \mu \mathrm{~s} / 1024$ points or about 1 ns apart. Therefore, when plotted versus reciprocal time, or frequency, they range from $1 \mathrm{MHz}(1 / 1 \mu \mathrm{~s})$ to about $1 \mathrm{GHz}(1 / 1 \mathrm{~ns})$. As shown in fig. 6.5, then, the signals of fig. 6.4 are contained in the frequency interval from $10^{6} \mathrm{~b} / \mathrm{s}$ to $10^{9} \mathrm{~b} / \mathrm{s}$.

Since it is economically unfeasible to calculate these step responses beyond the order of $1 \mu s$, (or equivalently, below $1 \mathrm{Mb} / \mathrm{s}$ ), due to computer memory and speed limitations, a straight line approximation has been used for the step response voltages between $1 \mu \mathrm{~s}$ and 1 s or $1 \mathrm{Mb} / \mathrm{s}$ and $1 \mathrm{~b} / \mathrm{s}$. Because the actual "one" voltage lies above and the actual "zero" voltage lies below their corresponding straight lines in this time or frequency interval, this straight line approximation is a conservative or "worst case" approximation.

### 6.3 Code Dependence

From the shapes of the cable step responses of figs. 6.2 and 6.3 it is evident that intersymbol interference will occur at the receiver when a cable is placed in the system. Just as in the case of an RC exponential response, the positive modeled cable step response reaches $V_{f}$ only as $t \rightarrow \infty$. Therefore, the voltage sampled at the receiver at the end of a bit period will depend not only on whether a "one" or a "zero" was sent in that period but also on the particular sequence of "ones" and "zeros" sent for all time before the bit period in question. In other words, the presence of this intersymbol interference, caused by the infinitely long cable step response "tails", implies that the system is code dependent.

One approach to the analysis of this code dependence on system operation would be to analyze separately the effects of all possible code sequences. However, a much simpler method exists which is to find the worst case and best case code sequences as lower and upper bounds for correct system operation.

The worst case is exactly the one referred to in the last section; an infinite string of "zeros" for all $t<0$, a transition to "one" at $t=0$ and an infinite string of "ones" for $t>0$. In other words it is represented by the (positive or negative) cable step response. If even a single "one" had occurred at some time $\mathrm{t}<0$, then the voltage at the receiver would not be exactly 0 V at $\mathrm{t}=0$. Rather, it would be slightly greater than 0 V due to the fact that the decay back to 0 V , even from a single "one" only occurs as $t \rightarrow \infty$. Therefore, if one or more "ones" occurred before $t=0$ and if a "one" is transmitted at $t=0$ as described above, then the step response commencing at $t=0$ will begin from some voltage greater than 0 V . The received voltage will thus have a "head start" towards $\mathrm{V}_{\mathrm{TH}}$ and will rise to that value sooner than in the case of all "zeros" preceding $t=0$. Of all possible code sequences preceding $t=0$, then, the one consisting of all "zeros" will yield the worst case because it is the only one that forces the transition to a "one" to begin at 0 V . The same argument may be used to show that the worst case for a "one"-to "zero" transition is the negative cable step response.

The best case is the sequence "one"- "zero"-"one"-"zero" occurring over all time. Fig. 6.6 (a through f) shows that this is true. Figure 6.6a is a plot of the transmitter signal for 64 bits of a "one"-"zero"-"one"-"zero" code sequence. (For all $t<0$ the transmitter output is 0 V.) Figure 6.6 b is the time domain convolution of this transmitter signal with the impulse response of 1050 feet of cable $I$, that is, figure 6.6 b is a plot of the receiver signal that results from the transmitter signal passing through the above cable at a $4 \mathrm{Mb} / \mathrm{s}$ rate. The plots shown in fig. 6.6 c through f are of the receiver signal that would be observed if the transmitter bit rate were increased to $8,16,32$ and $64 \mathrm{Mb} / \mathrm{s}$ respectively.

All of these responses may be thought of as transient responses starting at 0 V and eventually oscillating symmetrically about $\mathrm{V}_{\mathrm{TH}}$ as $\mathrm{t} \rightarrow \infty$. The worst case mentioned above is graphically illustrated by observing the first "zero" to "one" transition occurring at $t=0$. For the responses shown in fig. $6.6 \mathrm{~b}, \mathrm{c}$ and d , corresponding to 4,8 and $16 \mathrm{Mb} / \mathrm{s}$, respectively, the first "one" rises to a voltage greater than $V_{T H}$ thus implying that no errors will be made. For the 32 and $64 \mathrm{Mb} / \mathrm{s}$ responses, however, the first "one" does not reach $\mathrm{V}_{\mathrm{TH}}$; an error in detection of this "one" is certain. This tends to confirm the conclusion drawn from fig. 6.5 that the cutoff bit rate for proper system operation is about $20 \mathrm{Mb} / \mathrm{s}$.

The fact that the "one"-"zero"-"one"-"zero" sequence over all time is the best case may be deduced from the five responses shown in fig. 6.6 by observing that, regardless of the bit rate, all of the responses end up oscillating symmetrically about $V_{T H}$ as $t \rightarrow \infty$. Therefore, even for very high bit rates, when the receiver response signal reaches steady state the detector will properly discriminate between a "zero" and a "one".

Neither of these code sequences, the single step over all time or the square wave over all time, contain any information, of course, but they do represent the code boundaries. That is, for a given bit rate, they represent the minimum and maximum number of transitions allowable per unit of time. Therefore, any arbitrary "real" code sequence will have a receiver response somewhere between these worst and best case responses. Consequently, the worst case plot shown in fig. 6.5 may be used to prudently estimate the maximum bit rate that may be transmitted over a given type and length of cable.

### 6.4 Bit-Error-Rate Degradation

In the previous sections in this chapter it was shown that the overlay plot of the positive and negative cable step responses versus bit rate may be used to estimate the maximum bit rate that may be transmitted over a given type and length of cable under ideal conditions. Quite often, though, there is noise present in communications systems such that even if the system were operating at a bit rate less than the cutoff bit rate, there is a nonzero probability of error simply due to the noise. This error probability is usually quantitatively characterized by the parameter "bit error rate" (BER) which is the probability that an error will be made in the detection of any given bit. A BER of $10^{-6}$, for example, means that the probability of misinterpreting any given bit is one in one million, or, stated another way, it can be expected on the average that for every one million bits sent, one will be misinterpreted.

The goal of this section is to develop a method for determining the change in BER of an idealized communication system by the addition of a given type and length of cable. The assumed communication system is the same as that defined in section 6.1 and shown in fig. 6.1 but with the noise source switch
now closed. It will be further assumed that the noise source voltage output is Gaussian white noise; that is, uncorrelated noise with a Gaussian density function given as,

$$
\begin{equation*}
p(v)=\frac{1}{\sqrt{2 \pi \sigma^{2}}} e^{-(v-\mu)^{2} / 2 \sigma^{2}} \tag{6-2}
\end{equation*}
$$

where $\mu$ is the noise mean, assumed to be zero, and $\sigma$ is the noise standard deviation or rms value.
In order to determine the BER degradation due to the insertion of a cable into the system it is first necessary to calculate the BER of the system with the transmitter connected directly to the receiver. In this case a "one" is represented by a transmitter output of $V$ volts and a "zero" by the 0 volts. As before, the receiver threshold voltage is set to the midpoint of the received signal range, in this case $\mathrm{V} / 2$ volts. To this signal, then, is added a Gaussian noise voltage with zero mean and rms value $\sigma$.

If a "zero" is assumed to be sent then the probability of an error in detection is simply the probability that the received voltage (signal plus noise) will exceed v/2 volts. From the assumed Gaussian density function this is,

$$
\begin{equation*}
P(v>v / 2)=\int_{V / 2}^{\infty} \frac{1}{\sqrt{2 \pi \sigma^{2}}} e^{-v^{2} / 2 \sigma^{2}} d v \tag{6-3}
\end{equation*}
$$

Likewise, if a "one" is assumed to be sent then the probability of detection error is the probability that the received voltage is less than $V / 2$ and is given as

$$
\begin{equation*}
P(v<v / 2)=\int_{-\infty}^{v / 2} \frac{1}{\sqrt{2 \pi \sigma^{2}}} e^{-(v-v)^{2} / 2 \sigma^{2}} d v . \tag{6-4}
\end{equation*}
$$

Due to the symmetry of the Gaussian density function the probabilities of mistaking a "one" for a "zero" (eqn. 6-3) or a "zero" for a "one" (eqn. 6-4) are the same so either equation may be used to calculate the probability of error per bit regardless of whether a "one" or a "zero" was sent.

The actual calculation of the error probability from eqn (6-3) or (6-4) is quite difficult and fortunately unnecessary. Tables exist for the error function ${ }^{[11]}$, defined as,

$$
\begin{equation*}
\operatorname{erf}(x)=\frac{2}{\sqrt{\pi}} \int_{0}^{x} e^{-y^{2}} d y \tag{6-5}
\end{equation*}
$$

and it is easily shown [12] that the error probability above may be expressed in terms of the error function as

$$
\begin{equation*}
P(E)=\frac{1}{2}\left[1-\operatorname{erf}\left(\frac{V}{2 \sqrt{2} \sigma}\right)\right] \tag{6-6}
\end{equation*}
$$

Furthermore, if the voltage peak signal to rms noise ratio is defined as

$$
\begin{equation*}
\operatorname{VSNR}=\frac{\mathrm{V}}{\sigma} \tag{6-7}
\end{equation*}
$$

then the error probability per bit, or equivalently, the BER may be expressed as a function of VSNR as,

$$
\begin{equation*}
P(E)=\frac{1}{2}\left[1-\operatorname{erf}\left(\frac{\mathrm{VSNR}}{2 \sqrt{2}}\right)\right] \tag{6-8}
\end{equation*}
$$

A computer-generated plot of BER versus VSNR is given in fig. 6.7a through e.
As an example of using this plot assume that the transmitter output for a "zero" is 0 volts, a "one" is 24 volts and the rms noise value is 2 volts. Then, from (6-7) the VSNR is 12. From the plot in fig. 6.7c it is observed that a VSNR of 12 yields a BER of $10^{-9}$.

If the transmitter-receiver pair VSNR is known, it becomes possible to determine the new BER that would be observed by inserting of a given type and length of cable. This may be accomplished by use of the overlay step response versus bit rate plot of fig. 6.5 in conjunction with the BER versus VSNR plot of fig. 6.7 in the following manner. For a desired bit rate, the new peak signal voltage is obtained from fig. 6.5 as

$$
\begin{equation*}
v_{p}(B R)=v_{1}(B R)-v_{0}(B R) \tag{6-9}
\end{equation*}
$$

That is, the new peak signal voltage is the difference between the "one" voltage and the "zero" voltage that are present at the receiver at the end of a bit period for the assumed bit rate. From fig. 6.5, for example, $v_{p}$ for an assumed bit rate of $1 \mathrm{Mb} / \mathrm{s}$ is $0.86 \mathrm{~V}-0.07 \mathrm{~V}=0.79 \mathrm{~V}$. Then, from eqn. ( $6-7$ ), a new VSNR is obtained as

$$
\begin{equation*}
\operatorname{VSNR}_{\text {System }}+\text { Cable }=\frac{\mathrm{v}_{\mathrm{p}}}{\sigma}=\frac{0.79 \mathrm{~V}}{\sigma}=0.79 \mathrm{~V} \times \text { VSNR }_{\text {System }} \tag{6-10}
\end{equation*}
$$

Continuing the above example, if the transmitter receiver pair VSNR were 12, implying a system BER of $10^{-9}$, and then 1050 ft . of cable $I$ were placed in the system operating at $1 \mathrm{Mb} / \mathrm{s}$, the new system VSNR would be $0.79 \times 12$ or about 9.5. From fig. 6.7 b a VSNR of 9.5 yields a new BER of about $10^{-6}$.

These two graphs, figs. 6.5 and $6.7 a-e$, may be used for a variety of system design purposes. As another example suppose that it is desired to know the maximum permissible bit rate for the system in the previous example such that the new BER (after insertion of the cable) will not exceed $10^{-7}$. From fig. 6.7b this BER implies a minimum allowable VSNR of about 10.4. Then, from eqn. (6-10)

$$
\begin{equation*}
\mathrm{v}_{\mathrm{p}}=\mathrm{v} \frac{\mathrm{VSNR}_{\text {System }}+\text { Cable }}{\mathrm{VSNR}_{\text {System }}} \tag{6-11}
\end{equation*}
$$

or

$$
\mathrm{v}_{\mathrm{p}}=\mathrm{v} \quad \frac{10.4}{12}=0.87 \mathrm{~V}
$$

where in fig. 6.5 it is seen that this difference voltage $v_{p}$ occurs for a bit rate of $200 \mathrm{~b} / \mathrm{s}$.
However, rather than using the difference voltage of eqn. 6-9, one can proceed using a single curve, e.g., the "one" curve, fig. 6.2. Due to the symmetry of the negative and positive step response, it is evident that, for any bit rate

$$
\begin{equation*}
v_{0}(B R)=v_{f}-v_{1}(B R) \tag{6-12}
\end{equation*}
$$

$$
\begin{align*}
& \text { and substitution of this expression for } v_{0}(B R) \text { into (6-9) yields } \\
& \qquad v_{p}(B R)=2 v_{1}(B R)-v_{f}  \tag{6-13}\\
& \text { or }  \tag{6-14}\\
& v_{1}(B R)=\frac{1}{2}\left[v_{p}(B R)+v_{f}\right] .
\end{align*}
$$

For the example,

$$
\begin{aligned}
\mathrm{v}_{1}(\mathrm{BR}) & =\frac{1}{2}[0.87 \mathrm{~V}+0.926 \mathrm{~V}]=0.898 \mathrm{~V} \\
& \approx 0.9
\end{aligned}
$$

Observation of the point where the $v_{1}(B R)$ curve (the "zero"-to-"one" step response curve) passes through the 0.9 V line again yields a bit rate of approximately $200 \mathrm{~b} / \mathrm{s}$. The same answer has been obtained by reading the value of only one curve rather than by reading the difference between two curves.
CABLE


Figure 6.3. "One"-to-"zero" step response of 320.04 meters (1050 ft.) of cable I.



$$
\begin{aligned}
& \text { Figure 6.5. Superposition of positive and negative step responses of } 320.04 \text { meters } \\
& \text { ( } 1050 \mathrm{ft} .) \text { of cable I plotted as a function of bit rate. }
\end{aligned}
$$




Figure 6.6d. Received signal response for transmitter bit rate of $16 \mathrm{Mb} / \mathrm{s}$ through 320.04 meters ( 1050 ft. ) of cable I.






11.0
Voltage signal－－to－noise ratio
Figure 6．7c．$P(E)$ versus VSNR for VSNR from 10.0 to 14.0.
子！̣q xəd גоメコə ¥0 K7！t！qqeqoxd
7！̣q גəd エoxメə ¥0 K孔！T！qeqoxd

Voltage signal-to-noise ratio
Figure 6.7e. $P(E)$ versus VSNR for VSNR from 14.0 to 16.0 .

## 7. CABLE TESTING USING TIME DOMAIN REFLECTOMETRY

In this chapter the general principles of time domain reflection and transmission are developed and then are applied to cable testing. It is shown how the cable sending-end step responses are employed to calculate (simulate) the reference echoes (TDR responses) which characterize an undamaged cable. Departures from the reference response waveforms indicate uniform degradation of the cable and physical damage, or some combination of the two factors.

### 7.1 The General Theory of Time Domain Reflection and Transmission

In the classical frequency domain analysis of transmission lines containing reflections, the assumption is implicit that the sinusoidal signals are steady state signals, i.e., they have been present since the time minus infinity. In time domain measurements of transient or pulse phenomena, such is not the case. Prior to some time, say $t=0$, there was no signal present. When the pulse signal does enter the transmission line, it propagates along the line at a finite velocity arriving at observation points displaced from the input later in time. When this propagating signal encounters a discontinuity, reflected and transmitted signals are created. An observer remote from the point of reflection will not observe the reflected signal until still later in time due to the time isolation (delay) inherent in the transmission line. For people working with time domain measurements and instruments it is very important to comprehend this time isolation of reflections. As an aid to developing an understanding of this concept this section develops the analysis of a discretely loaded transmission line in terms of transmitted and reflected pulse trains. A reflection-transmission (RT) diagram is used. The resultant transmission line equations are then shown to be identical to those derived in the classical frequency domain manner.

The starting point for this discussion is based upon fig. 2.2 and the general equations (2-3) and (2-4). First, equation (2-4) is expressed in terms of $\mathrm{E}_{\mathrm{g}}(\mathrm{s})$ rather than $\mathrm{I}_{\mathrm{s}}$ (s) through the relation

$$
\begin{equation*}
I_{s}(s)=\frac{E_{g}(s)-E_{s}(s)}{Z_{g}(s)} \tag{7-1}
\end{equation*}
$$

with the result

$$
\begin{equation*}
\frac{E_{s}(s)}{E_{g}(s)}=\left[\frac{Z_{0}(s)}{Z_{g}(s)+Z_{o}(s)}\right]\left[\frac{1}{1-\rho_{S}(s) \rho_{r}(s) e^{-2 \ell \gamma(s)}}\right]\left[1+\rho_{r}(s) e^{-2 \ell \gamma(s)}\right] . \tag{7-2}
\end{equation*}
$$

For convenience, $(2-3)$ is repeated here,

$$
\begin{equation*}
\frac{E_{r}(s)}{E_{g}(s)}=2\left[\frac{Z_{o}(s)}{Z_{g}(s)+Z_{o}(s)}\right]\left[\frac{Z_{\ell}(s)}{Z_{\ell}(s)+Z_{o}(s)}\right]\left[\frac{e^{-\ell \gamma(s)}}{1-\rho_{s}(s) \rho_{r}(s) e^{-2 \ell \gamma(s)}}\right] \tag{7-3}
\end{equation*}
$$

$\rho_{s}(s)$ and $\rho_{r}(s)$ are the reflection coefficients (2-5) produced by the sending and receiving end impedances, $Z_{g}(s)$ and $Z_{\ell}(s)$, respectively. \& is the transmission line length while $\gamma(s)$ and $Z_{o}(s)$ are the line propage*.ion function and characteristic impedance, respectively.

$$
\begin{align*}
& Z_{0}(s)=\left[\frac{Z(s)}{Y(s)}\right]^{1 / 2}  \tag{2-1}\\
& \gamma(s)=[Z(s) Y(s)]^{1 / 2} \tag{2-2}
\end{align*}
$$

where $Z(s)$ and $Y(s)$ are the transmission line equivalent circuit parameters: series impedance/unit length and shunt admittance/unit length. In general, for wave transmission on a transmission line $\gamma(s)$ is such that the limit

$$
\lim _{s \mid \rightarrow \infty}\left[\frac{r(s)}{s}\right]
$$

is equal to the time delay per unit length, T, i.e.,

$$
\begin{equation*}
\mathrm{T}=\lim _{|\mathrm{s}| \rightarrow \infty}\left[\frac{\gamma(\mathrm{s})}{\mathrm{s}}\right] \tag{7-4}
\end{equation*}
$$

For uniform transmission lines, the time delay/unit length is expressed in terms of the inductance and capacitance per unit length,

$$
\begin{equation*}
\mathrm{T}=\sqrt{\mathrm{LC}} \tag{7-5}
\end{equation*}
$$

Also, $\gamma(s)$ can be written in the form

$$
\begin{equation*}
\gamma(s)=s T+\gamma^{\prime}(s) \tag{7-6}
\end{equation*}
$$

where sT is an all-pass function (linear phase function) and $\gamma^{\prime}$ (s) is the minimum phase function (see section 2.4 for some discussion of minimum phase). For $\gamma^{\prime}(s)$ to be a minimum phase function, the ratio $\gamma^{\prime}(s) / s$ vanishes as $|s|$ becomes infinitely large [2].

Now consider a step generator of impedance $Z_{o}(s)$ applied to the system described by (7-2) and (7-3), i.e., $E_{g}(s)=1 / s$ and $Z_{g}(s)=Z_{o}(s)$. Then $\rho_{s}(s)=0$, and $(7-2)$ and (7-3) reduce to

$$
\begin{equation*}
E_{s}(s)=\frac{1}{2 s}\left[1+\rho_{r}(s) e^{-2 \ell\left[s T+\gamma^{\prime}(s)\right]}\right] \tag{7-7}
\end{equation*}
$$

and

$$
\begin{equation*}
E_{r}(s)=\frac{1}{s}\left[\frac{Z_{\ell}(s)}{Z_{0}(s)+Z_{\ell}(s)}\right] e^{-\ell\left[s T+\gamma^{\prime}(s)\right]} \tag{7-8}
\end{equation*}
$$

Rewriting

$$
\begin{equation*}
E_{s}(s)=\frac{1}{2 s}+e^{-2 s \ell T}\left[\rho_{r}(s) e^{-2 \ell \gamma^{\prime}(s)}\right] \tag{7-9}
\end{equation*}
$$

and

$$
\begin{equation*}
E_{r}(s)=e^{-s \ell T}\left\{\frac{Z_{\ell}(s) e^{-\ell \gamma^{\prime}(s)}}{s\left[Z_{o}(s)+Z_{\ell}(s)\right]}\right\} \tag{7-10}
\end{equation*}
$$

Consideration of (7-9) and (7-10) leads to the following conclusions:

1. The input terminals of the line presents to the generator the characteristic impedance $Z_{o}(s)$ for the time interval $0<t<2 \ell T$; i.e., the load impedance is not seen at the generator terminals until the incident wave travels to the load and the subsequent reflection returns to the generator terminals.
2. The output waveform is zero for $\mathrm{t}<\ell \mathrm{T}$; also, its shape is controlled by the characteristic impedance to load impedance voltage divider and the minimum phase (distorting) part of the propagation function, $\gamma^{\prime}(s)$.

Generally, for arbitrary terminations, $E_{s}(s)$ and $E_{r}(s)$ each consist of an infinite sequence of terms similar to those in (7-9) and (7-10) and have their origins in the conclusions delineated above. By using the geometric series expansion

$$
\begin{equation*}
\left[1-\rho_{s}(s) \rho_{r}(s) e^{-2 \ell \gamma(s)}\right]^{-1}=\sum_{n=0}^{\infty}\left[\rho_{s}(s) \rho_{r}(s) e^{-2 \ell \gamma(s)}\right]^{n} ; \tag{7-11a}
\end{equation*}
$$

where

$$
\begin{equation*}
\left|\rho_{s}(s) \rho_{r}(s) e^{-2 \ell \gamma(s)}\right|^{2}<1 \tag{7-11b}
\end{equation*}
$$

in the second and third factors of (7-2) and (7-3), respectively, the transfer functions can be expressed as infinite sequences which represent impulse responses consisting of pulse trains. A note of caution should be entered here; for a lossless line the inequality condition in (7-11) may not be satisfied and hold for all values of $s$ in the complex plane. Under such conditions an alternate mathematical procedure must be used [13].

Putting (7-11) into (7-2) and (7-3) yields the sending end pulse train

$$
\begin{align*}
E_{s}(s) & =E_{g}(s)\left[\frac{Z_{0}(s)}{Z_{g}(s)+Z_{o}(s)}\right] \\
X & \sum_{n=0}^{\infty} \rho_{s}^{n}(s) \rho_{r}^{n}(s) e^{-2 n \ell \gamma(s)}+\rho_{s}^{n}(s) \rho_{r}^{(n+1)}(s) e^{-(n+1) 2 \ell \gamma(s)} \tag{7-12}
\end{align*}
$$

and the receiving end pulse train,

$$
\left.\left.\begin{array}{rl}
E_{r}(s)= & 2 E_{g}(s)
\end{array}\right] \frac{Z_{o}(s)}{Z_{g}(s)+Z_{o}(s)}\right]\left[\frac{Z_{\ell}(s)}{Z_{\ell}(s)+Z_{o}(s)}\right] .
$$

where (7-12) and (7-13) are both subject to the condition specified in (7-11b). The physical situation represented by the pulse trains (7-12) and (7-13) is lucidly shown by a reflection-transmission (RT) diagram, fig. 7.1. In fact, the RT diagram itself provides an intuitive (but exact) means for writing out the terms of each pulse train.

The RT diagram is a signal flow chart embodying the reflection and transmission properties of the terminal impedances, and the wave propagation property of the connecting transmission line. To interpret the diagram the following rules are used:

1. Passage from a generator of impedance $Z_{g}(s)$ through a node to a line segment represents multiplication by the transfer function $Z_{o}(s) /\left[Z_{g}(s)+Z_{o}(s)\right]$.
2. Passage along a line segment represents multiplication by the transfer function exp [-ly(s)].
3. Passage through a node to a load $Z_{\ell}(s)$ represents multiplication by the transfer function $2 Z_{\ell}(s) /\left[Z_{\ell}(s)+Z_{o}(s)\right]$.
4. Reflection from a node represents multiplication by the reflection coefficient

$$
p_{r}(s)=\left[Z_{\ell}(s)-Z_{o}(s)\right]\left[Z_{\ell}(s)+Z_{o}(s)\right]^{-1}
$$

Application of the rules to the diagram in fig. 7.1 obtains the following results:

$$
\begin{align*}
& E_{S}(s)=\left[E_{g} \frac{Z_{o}}{Z_{g}+Z_{o}}\right]+\left[E_{g} \frac{Z_{o}}{Z_{g}+Z_{o}} e^{-\ell \gamma_{o}} e^{-2 \gamma} \frac{2 Z_{g}}{Z_{g}+Z_{o}}\right] \\
& +\left[E_{g}(s) \frac{Z_{0}}{Z_{0}+Z_{g}} e^{-l \gamma_{\rho_{r}}} e^{-l \gamma_{\rho}} e^{-l \gamma_{\rho}} e^{-l \gamma} \frac{2 Z}{Z_{g}+Z_{o}}\right]+\ldots . \\
& =E_{g} \frac{Z_{o}}{Z_{g}+Z_{o}} \left\lvert\, 1+\frac{2 \rho_{r} Z_{g}}{Z_{g}+Z_{o}} e^{-2 \ell \gamma}+\frac{2 \rho^{2} \rho_{\rho}^{\rho} Z_{g}}{Z_{g}+Z_{o}} e^{-4 \ell \gamma}+\ldots .\right. \tag{7-14}
\end{align*}
$$

Since

$$
\rho_{i}=\frac{z_{i}-z_{o}}{z_{i}+z_{o}}
$$

the voltage divider may be written as

$$
\frac{Z_{i}}{Z_{i}+z_{0}}=\frac{1+p_{i}}{2}
$$

Putting ( $\left.1+\rho_{s}\right) / 2$ into (7-14) for the factor $Z_{g} /\left(Z_{g}+Z_{o}\right)$ obtains the result

$$
\begin{equation*}
E_{s}(s)=E_{g} \frac{Z_{o}}{Z_{g}+Z_{o}}\left[1+\rho_{r}\left(1+\rho_{s}\right) e^{-2 \ell \gamma}+\rho_{r}^{2} \rho_{s}\left(1+\rho_{s}\right) e^{-4 \ell \gamma}+\ldots \cdot\right] \tag{7-15}
\end{equation*}
$$

which is identical to that given by expansion of (7-12). Therefore, it has been demonstrated that the RT diagram allows one to write down the equation for $E_{s}(s)$ directly by tracing out the signal flow while invoking the rules $1-4$.

Similarly, the received voltage $E_{r}(s)$ is obtained from the $R T$ diagram as

$$
\begin{align*}
& E_{r}(s)=\left[E_{g} \frac{z_{o}}{z_{g}+z_{o}}\right]\left[\begin{array}{ll}
e^{-\ell \gamma} \frac{2 z_{\ell}}{z_{\ell}+z_{o}} \\
& \\
A B
\end{array}\right] \\
& +\left[E_{g} \frac{z_{o}}{z_{g}+z_{o}} e^{-\ell \gamma_{\rho_{r}}} e^{-\ell \gamma_{\rho_{s}}} e^{-\ell \gamma} \frac{2 z_{\ell}}{z_{\ell}+z_{o}}\right]+\ldots \cdot \\
& =2 E_{g} \frac{Z_{o}}{Z_{g}+Z_{o}} \frac{z_{\ell}}{Z_{\ell}+Z_{o}}\left[e^{-\ell \gamma}+\rho_{r^{\prime}} \rho e^{-3 \ell \gamma}+\ldots\right] \tag{7-16}
\end{align*}
$$

which is identical to that found by expanding (7-13).
In this report the transmission lines are lossy and are normally operated in a doubly mismatched condition due to the terminations being equal to the nominal high frequency characteristic impedance, $R_{0}$. Under such conditions, (7-15) and (7-16) become

$$
\begin{equation*}
E_{s}(s)=E_{g}(s) \frac{Z_{o}(s)}{R_{o}+Z_{o}(s)}\left[1+\rho_{r}\left(1+\rho_{s}\right) e^{-2 \ell r(s)}+\rho_{r} \rho_{s}\left(1+\rho_{s}\right) e^{-4^{\ell \gamma}(s)}+\ldots .\right] \tag{7-17}
\end{equation*}
$$

and

$$
\begin{equation*}
E_{r}(s)=2 E_{g}(s) \frac{Z_{o}(s) R_{o}}{\left[R_{o}+Z_{o}(s)\right]^{2}}\left[e^{-\ell \gamma(s)}+p^{2} e^{-3 \ell \gamma(s)}+\ldots\right] \tag{7-18}
\end{equation*}
$$

respectively, where

$$
\begin{equation*}
\rho=\rho_{r}=\rho_{s}=\frac{R_{0}-Z_{o}(s)}{R_{0}+Z_{o}(s)} \tag{7-19}
\end{equation*}
$$

Briefly, consider the received voltage $\mathrm{E}_{\mathrm{r}}(\mathrm{s})$. Inspection of (7-18) shows that the initial signal transfer function is due to the product of the RT input voltage transfer function, the exponential transmission factor, and the output RT voltage transfer function, i.e.,

$$
\begin{equation*}
\left.E_{r}(s)\right]_{0<t<3 \ell T}=2 E_{g}(s)\left[\frac{Z_{o}(s)}{Z_{o}(s)+R_{o}} \times e^{-\ell \gamma(s)} x \frac{R_{o}}{Z_{o}(s)+R_{o}}\right] \tag{7-20}
\end{equation*}
$$

This is in agreement with the RT diagram transmission/reflection rules. The condition $0<t<3 \ell T$ describes the time interval of the initial response, which is to say that the inverse Laplace transform of (7-20) represents the time domain response only over the interval $0<t<3 \ell T$. Refer to fig. $7-2$ and the brief disdiscussion in its caption.

$$
\begin{equation*}
e_{r}(t)=\varepsilon^{-1}\left[2 E_{g}(s) \frac{Z_{o}(s) R_{o}}{Z_{o}(s)+R_{o}} e^{-\ell \gamma(s)}\right] ; 0<t<3 \ell T . \tag{7-21}
\end{equation*}
$$

It is important to keep in mind that (7-21) is the exact response for $e_{r}(t)$ over the interval $0<t<3 \ell T$. What is unusual about this method of analysis is that even though a series expansion (7-11) has been used, a finite sequence of terms starting from the initial term is always an exact result out to some multiple of
the time lT. Contrast this with other series expansions where an infinite number of terms must be present for an exact result, e.g., Fourier Series, etc. The discussion now turns to the sending-end response.

The main purpose of this chapter is to develop the sending end response so that one would have a reference response or waveform for a given cable type; this response would then be used to evaluate the measured test response of samples of the given cable type. To conclude this section, consider (7-17). It is clear that for the time interval $0<t<2 \ell T$ the response is only dependent upon the transmission line through the characteristic impedance, $Z_{o}(s)$, and is given by the first term from the sequence (7-17)

$$
\begin{equation*}
\left.E_{s}(s)\right]=E_{g}(s) \frac{Z_{o}(s)}{R_{0}+Z_{o}(s)} \tag{7-22}
\end{equation*}
$$

In the next section this relation will be used to derive the reference responses. After the passage of $2 \ell T$ seconds, the second term in the sequence must be added to the first to provide the temporal response out to $4 \ell T$ seconds

$$
\begin{equation*}
\left.E_{s}(s)\right]_{0<t<4 \ell T}=E_{g}(s)\left[\frac{Z_{o}(s)}{R_{0}+Z_{o}(s)}+\frac{Z_{o}(s)}{R_{0}+Z_{o}(s)} \rho_{I}\left(1+\rho_{s}\right) e^{-2 \ell \gamma(s)}\right] \tag{7-23}
\end{equation*}
$$

Now consider the second term in detail; the term ( $1+p_{s}$ ) is just $2 R_{o} /\left[R_{o}+Z_{o}(s)\right]$ which is the $R T$ transfer function from the transmission line to the load $R_{o} ; \rho_{r}$ is the reflection coefficient of the receiving end load. $Z_{o}(s) /\left[R_{0}+Z_{o}(s)\right]$ is the $R T$ transfer function from the generator to the sending end, and exp.[-2ly(s)] represents transmission over a length $2 \ell$. Consequently, the second term has a physical picture which correlates with the RT diagram as follows:


In other words, the initial sending end voltage is transmitted to the load; the received voltage is then scaled by the reflection coefficient and then transmitted back to the sending end. Figure 7-3 qualitatively illustrates the time domain unit step response corresponding to (7-23). In the next section the time domain expressions for unit step excitation will be derived.

### 7.2 The Step Response of the Sending-End Voltage

The sending-end voltage $E_{S}(s)$ for $0 \leq t \leq 2 l t$ due to a unit step is obtained from (7-22) by letting

$$
\begin{equation*}
E_{g}(s)=\frac{1}{s} \tag{7-24}
\end{equation*}
$$

which obtains for the sending-end voltage

$$
\begin{equation*}
E_{s}(s)=\frac{1}{s} \frac{Z_{0}(s)}{R_{0}+Z_{0}(s)} ; \quad 0 \leq t \leq 2 \ell T \tag{7-25}
\end{equation*}
$$

The time domain sending-end voltage is given by

$$
\begin{equation*}
e_{s}(t)=\mathcal{L}^{-1}\left\{\frac{1}{s} \frac{Z_{0}(s)}{R_{o}+Z_{o}(s)}\right\} \quad 0 \leq t \leq 2 \ell T \tag{7-26}
\end{equation*}
$$

This expression is also valid for all time $t$ when the receiving-end is terminated in the characteristic impedance, $Z_{o}(s)$, that is; when there are no reflections from the receiving-end, the line appears to be infinitely long. For a finite length, $\ell, e_{s}(t)$ is only valid for $0 \leq t \leq 2 \ell T$. In terms of a circuit model, $e_{s}(t)$ represents a voltage division of the generator voltage, $e_{g}(t)$, fig. 7-4A. Qualitatively, the response is shown in fig. 7.3. If there were no reflection from the load (receiving-end), the waveform would continue to rise monotonically to unity, fig. 7.4B. From the final value theorem of the Laplace transformation, the final value of (7-26) is

$$
\begin{equation*}
e_{s}(\infty)=\lim _{|s| \rightarrow 0} s\left[\frac{1}{s} \frac{Z_{o}(s)}{R_{0}+Z_{o}(s)}\right]=1 \tag{7-27}
\end{equation*}
$$

Note that (7-26) is independent of $\ell$; that is to say, the step response of any length of cable will follow (7-26) until the first reflection arrives at the sending-end ( $t=2 \ell T$ ).

Expressing (7-25) in terms of the model parameters gives

$$
\begin{align*}
& E_{s}(s)=\frac{1}{s} \frac{\left[\frac{R+s L+K s^{m}}{s C}\right]^{1 / 2}}{m]^{1 / 2}} \\
& R_{0}+\left[\frac{R+s L+K s^{m}}{s C}\right] \\
& =\frac{1}{s} \frac{R_{o}\left[1+\frac{R+K s^{m}}{s L}\right]^{1 / 2}}{R_{o}+R_{o}\left[1+\frac{R+K s^{m}}{s L}\right]^{1 / 2}} \\
& E_{s}(s)=\frac{1}{s} \frac{\left[1+\frac{R+K s^{m}}{s L}\right]^{1 / 2}}{1 / 2} .  \tag{7-28}\\
& 1+\left[1+\frac{R+K s^{m}}{s L}\right]
\end{align*}
$$

For large $s$ (small t) (7-28) can be expanded as the ratio of two power series,

$$
\begin{equation*}
E_{s}(s)=\frac{1}{s} \frac{1+\frac{v}{2}-\frac{v^{2}}{2 \cdot 4}+\frac{3 v^{3}}{2 \cdot 4 \cdot 6}-\frac{3.5 v^{4}}{2 \cdot 4 \cdot 6 \cdot 8}+\cdots}{2+\frac{v}{2}-\frac{v^{2}}{2 \cdot 4}+\frac{3 \cdot v^{3}}{2 \cdot 4 \cdot 6}-\frac{3.5 v^{4}}{2 \cdot 4 \cdot 6 \cdot 8}+\cdots} \tag{7-29}
\end{equation*}
$$

where

$$
\begin{equation*}
v=\frac{R+K s^{m}}{s L} \tag{7-30}
\end{equation*}
$$

and

$$
\begin{equation*}
|v|=\left|\frac{R+K s^{m}}{s L}\right|<1 \tag{7-31}
\end{equation*}
$$

If the condition (7-31) is violated, the expansion (7-29) is not valid. By dividing the denominator of (7-29) into its numerator a single series in $v$ can be obtained. This is easily done using the following method. If the product of two polynominals in $v$ is given by

$$
\mathrm{Y}=\mathrm{HX}
$$

then the k -th coefficient of Y is given by

$$
\begin{equation*}
y(k)=\sum_{i=1}^{k_{1}} h(i) x(k+1-i) . \tag{7-32}
\end{equation*}
$$

For example,

$$
\begin{aligned}
& y(1)=h(1) x(1) \\
& y(2)=h(1) x(2)+h(2) x(1)
\end{aligned}
$$

and

$$
Y=h(1)+h(2) v+h(3) v^{2}+\cdots
$$

By writing out a few terms in the product, and starting with $x(1) \neq 0$, it is possible to solve for $h(k)$, i.e., the quotient of $Y / X$. The result is

$$
\begin{equation*}
h(k)=\frac{1}{x(1)}\left[y(k)-\sum_{i=1}^{k-1} h(i) x(k+1-1)\right] \tag{7-33}
\end{equation*}
$$

where

$$
\begin{aligned}
h(1) & =\frac{y(1)}{x(1)} \\
h(2) & =\frac{1}{x(1)}[Y(2)-h(1) x(2)] \\
& \vdots
\end{aligned}
$$

Consequently (7-29) can be replaced by

$$
\begin{equation*}
E_{s}(s)=\frac{1}{2 s}+\frac{1}{8 s} v-\frac{1}{16 s} v^{2}+\frac{5}{128 s} v^{3}-\frac{14}{512 s} v^{4}+\cdots \tag{7-34}
\end{equation*}
$$

The series (7-34) is also convergent for

$$
\begin{equation*}
|v(s)|=\left|\frac{R+K s^{m}}{s L}\right|<1 . \tag{7-35}
\end{equation*}
$$

Upon substituting $v(s)$ into the series (7-34) there results a series having terms of the form $\mathrm{As}^{-\mathrm{q}}$ and $\mathrm{Bs}^{-r}$ where q is an integer $1,2,3 \ldots$, and $r$ is not an integer. The time domain response is obtained by taking the inverse Laplace transform of each term using the appropriate transform pair

$$
\begin{align*}
& \frac{A}{s^{q}} \longleftrightarrow A \frac{t^{q-1}}{q-1} u(t)  \tag{7-36}\\
& -\frac{B}{s^{r}} \leftrightarrow \frac{B t^{r-1}}{\Gamma(r)} u(t) \tag{7-37}
\end{align*}
$$

where $u(t)$ is the unit step and $\Gamma(r)$ the gamma function. The non integers $r$ result from the fractional power $m, 0<m<1$, which appears in the high frequency loss term of the model, $\mathrm{Ks}^{\mathrm{m}}$. The resultant time domain response is

$$
\begin{equation*}
e_{S}(t)=\frac{1}{2} u(t)+\frac{1}{8} w_{1}(t)-\frac{1}{16} w_{2}(t)+\frac{5}{128} w_{3}(t)-\frac{7}{256} w_{4}(t)+\cdots \tag{7-38}
\end{equation*}
$$

or in terms of the coefficient $b_{n}$,

$$
\begin{equation*}
e_{s}(t)=\frac{1}{2} u(t)+\sum_{n=1}^{\infty} b_{n} w_{n}(t) \tag{7-39}
\end{equation*}
$$

where

$$
\begin{align*}
& w_{n}(t)=\sum_{i=0}^{n} \frac{a_{n i} R^{n-i} K^{i}}{L^{n} \Gamma[(n+1)-i m]} t^{n-i m} u(t)  \tag{7-40}\\
& a_{n i}=\frac{n!}{i!(n-i)!}  \tag{7-41}\\
& b_{n}=\frac{\prod_{i=1}^{n}(2 i-1)}{(2 n+2) \prod_{i=1}^{n}(2 i)}(-1)^{n+1}  \tag{7-42}\\
& n!=\Gamma(n+1)  \tag{7-43}\\
& \Gamma(n+1)=n \Gamma(n) . \tag{7-44}
\end{align*}
$$

For examples of the expansion of $w_{n}(t),(7-40), w_{1}(t)$ through $w_{4}(t)$ are given below.

$$
\begin{align*}
& w_{1}(t)=\frac{R}{L \Gamma(2)} t u(t)+\frac{K}{L \Gamma(2-m)} t^{1-m} u(t)  \tag{7-45}\\
& w_{2}(t)= \frac{R^{2}}{L^{2} \Gamma(3)} t^{2} u(t)+\frac{2 R K}{L^{2} \Gamma(3 m)} t^{2-m} u(t)+\frac{R^{2}}{L^{2} \Gamma(3-2 m)} t^{2-m} u(t)  \tag{7-46}\\
& w_{3}(t)= \frac{R^{3}}{L^{3} \Gamma(4)} t^{3} u(t)+\frac{3 R^{2} K}{L^{3} \Gamma(4-m)} t^{3-m} u(t)+\frac{3 R K^{2}}{L^{3} \Gamma(4-2 m)} t^{3-2 m} u(t) \\
&+\frac{K^{3}}{L^{3} \Gamma(4-3 m)} t^{3-3 m} u(t)  \tag{7-47}\\
& w_{4}(t)=\frac{R^{4}}{L^{4} \Gamma(5)} t^{4} u(t)+\frac{4 R^{3} K}{L^{4} \Gamma(5-m)} t^{4-m} u(t)+\frac{L^{4}}{L^{4} \Gamma(5-2 m)} t^{4-2 m} u(t) \\
&+\frac{4 R^{2} K^{2}}{L^{4} \Gamma(5-3 m)} t^{4-3 m} u(t)+\frac{K^{4}}{L^{4} \Gamma(5-4 m)} t^{4-4 m} u(t) \tag{7-48}
\end{align*}
$$

For a given cable, certain precautions must be taken before using the time domain series (7-38) to calculate the sending-end voltage $e_{s}(t)$. First of all, the condition (7-35) must be satisfied for all frequencies comprising the sending-end step response voltage $e_{s}(t)$. As can be seen from (7-35), there is a lower limit on $|s|$, i.e., a lower frequency limit, $f_{0}$. Letting $s=j \omega$ in (7-35) obtains

$$
\begin{equation*}
\left|\frac{R / L}{j \omega}+\frac{(K / L)}{(j \omega)^{1-m}}\right|<1 \tag{7-49}
\end{equation*}
$$

which can be expressed as

$$
\begin{equation*}
\frac{\mathrm{K}^{2}}{\omega^{2(1-\mathrm{m})} \mathrm{L}^{2}}+\frac{\mathrm{R}^{2}}{\omega^{2} \mathrm{~L}^{2}}+\frac{2 \mathrm{RK}}{\omega^{2-m} \mathrm{~L}^{2}} \sin \frac{(1-\mathrm{m}) \pi}{2} \quad 1 / 2<1 \tag{7-50}
\end{equation*}
$$

The frequency, $f=\omega / 2 \pi$, at which the left hand side of (7-50) just equals unity is defined as $f_{0}$. For frequencies equal to and less than $f_{0}$, the series (7-34) will not converge to $E_{s}(s)$. Consequently, the corresponding time domain series (7-38) will diverge for time greater than some valute $T_{1}$, which should be of the order of the reciprocal of $f_{o}$, i.e.,

$$
T_{0}=1 / f_{0}
$$

$T_{1}$ can not be exactly specified due to the uncertainty properties of the Fourier (and Laplace)
transformation which prohibits a one to one correspondence between points in the time and frequency domain.

To demonstrate the application of the time domain step response, equation (7-38), and also its convergence properties, the sending-end step response of Cable I has been simulated using (7-38); the cable parameters are given in section 5.1. Figures $7-5$ and $7-6$ show the responses for the time windows of 5 and 10 microseconds, respectively. Note the initial jump to 0.5 and the subsequent slow rise thereafter increasing only by 10 to $12 \%$ of the final value (1.0). For each response the largest error occurs at the last point; they are $1.5 \times 10^{-4}$ volts ( $1.5 \times 10^{-2} \%$ ) and $3.3 \times 10^{-3}$ volts ( $0.33 \%$ ) for figs. 7.5 and 7.6 , respectively. The errors occur due to the use of a finite number of terms of the series $e_{s}(t)$, (7-38). For figs. 7.5 and 7.6 the truncated series for (7-38) consisted of the first term $0.5 u(t)$ plus the weighted values of $w_{1}(t)$ through $w_{6}(t)$, a total of 28 terms. The error was computed by taking the magnitude of the difference between results for the truncated series and a second truncated series containing one more weighted value of $w_{i}(t)$. The second truncated series contained $w_{8}(t)$ and thus contained a total of 36 terms.

Turning now to the divergent properties of the series, it is known that (7-34) will not converge for values of $|s|$ along the real frequency axis, i.e., $|\omega|$, when $\omega$ is equal to or less than $2 \pi \times 28 \times 10^{3}$, $f_{0}=28 \mathrm{kHz},(7-50)$. Consequently, in the neighborhood of $t=T_{o}=1 / f_{o}=35.7$ microseconds, the time domain response ( $7-38$ ) should diverge even when an infinite number of its terms are used. Truncation of the series should lead to a divergent result well before 35.7 microseconds; that such is the case is shown in fig. 7-7.

Figure 7-7 presents the step responses for the time interval of 40 microseconds as a function of the number of terms of (7-38). The label 10 corresponds to 10 terms comprised of $0.5 u(t)$ plus the algebraic sum of the weighted values of $w_{1}(t)$ through $w_{3}(t)$. Similarly, 15 corresponds to 15 terms obtained from the algebraic sum of the weighted terms through $w_{4}(t)$. The largest number of terms, 36 yields the most accurate result. For example, at $t=10$ microseconds the responses for 28 and 36 terms are virtually identical while those for 21,15 , and 10 terms show increasing error, respectively. Also, the divergence increases as $t$ approaches the neighborhood of $T_{0}, 35.7$ microseconds. These curves show that the divergent properties of the series (7-38) and its truncation to 28 or 36 terms have negligible effect on the computation of $e_{s}(t)$ out to 10 microseconds (no greater than $0.33 \%$ error at 10 ns ).

For the cables considered here, the maximum physical length is typically about 300 m ( 1000 ft ). Accordingly, for a typical relative dielectric constant of about 2.5 , the maximum cable delay to be encountered would be about 1.69 microseconds ( $\sqrt{2.5} \times 320 / 3 \times 10^{8}$ ) and the corresponding $2 \ell T$ value would be 3.38 microseconds ( $2 \times 1.69 \times 10^{-6}$ ). In the particular case of Cable $I$, $\ell T$ is equal to $1.6 \times 10^{-6}$ ( $320 \times \sqrt{\mathrm{LC}}$ ) and $2 \ell \mathrm{~T}$ is 3.2 microseconds. Accordingly, the observed response would depart from (7-38) at 3 ns in the manner shown in fig. 7-3; however, figs. 7-5 and 7-6 do accurately describe the observed response out to 3.2 ns , i.e., for the time interval before the receiving-end reflection returns to the sending-end.

### 7.3 A Practical Method for Measuring the Step Response of the Sending-End Voltage

When the generator impedance $R_{g}$ is not equal to the high frequency characteristic impedance $R_{0}$ of the cable under test, (7-28) is not valid. Specifically, in the first form of (7-28) note the $R_{0}$ term; for a generator of impedance $R_{g}$, the $R_{0}$ term has to be replaced by $R_{g}$. Consequently, for a generator impedance $R_{g}$ all equations based on (7-28) are also invalid.

However, all is not lost because the principal result of Section 7.2 , es (t), i.e., eq. ( $7-38$ ), transforms as follows:

$$
\begin{equation*}
\left.e_{s}(t)\right]_{R_{g}}=\frac{2 R_{o}}{R_{o}+R_{g}} e_{s}(t) \tag{7-51}
\end{equation*}
$$

The derivation of this result proceeds in the following manner. Consider fig. 7-4A in which $R_{0}$ is replaced by $R_{g}$. Then, the voltage across $Z_{o}(s)$ is given by

$$
\begin{equation*}
\left.e_{s}(t)\right]_{R_{g}}=\frac{1}{s} \frac{Z_{o}(s)}{R_{g}+Z_{o}(s)} \tag{7-52}
\end{equation*}
$$

Upon expressing $Z_{o}(s)$ in terms of the cable parameters obtains

$$
\begin{align*}
\left.e_{s}(t)\right]_{g} & =\frac{1}{s} \frac{\left[\frac{R+s L+K s^{m}}{s L}\right]^{1 / 2}}{R_{g}+\left[\frac{R+s L+K s^{m}}{s L}\right]^{1 / 2}} \\
& =\frac{1}{s} \frac{R_{o}\left[1+\frac{R+k s^{m}}{s L}\right]^{1 / 2}}{R_{g}+\left[1+\frac{R+K s^{m}}{s L}\right]^{1 / 2}} \\
& =\frac{1}{s} \frac{\left[1+\frac{R+K s^{m}}{s L}\right]^{1 / 2}}{R_{g}+\left[1+\frac{R+K s^{m}}{s L}\right]^{1 / 2}} \tag{7-53}
\end{align*}
$$

where $\bar{R}_{g}$ denotes $R_{g} / R_{0}$. Expanding the radical terms in a power series as was done in (7-28) yields

$$
\begin{equation*}
\left.e_{s}(t)\right]_{R_{g}}=\frac{1+\frac{v}{2}-\frac{v^{2}}{2 \cdot 4}+\frac{3 v^{3}}{2 \cdot 4 \cdot 6}-\frac{3 \cdot 5 \cdot v^{4}}{2 \cdot 4 \cdot 6 \cdot 8}+\cdots}{\left(1+\bar{R}_{g}\right)+\frac{v}{2}-\frac{v^{2}}{2 \cdot 4}+\frac{3 v^{3}}{2 \cdot 4 \cdot 6}-\frac{3 \cdot 5 \cdot v^{4}}{2 \cdot 4 \cdot 6 \cdot 8}+\cdots} \tag{7-54}
\end{equation*}
$$

Because the quotient of polynomials, (7-54), can be expressed as

$$
H=\frac{Y}{X}=\frac{1+a_{1} v+a_{2} v^{2}+\cdots}{\left(1+\bar{R}_{g}\right)+a_{1} v+a_{2} v},
$$

the $k$-th coefficient of $H$ is given by

$$
\begin{equation*}
h(k)=\frac{1}{x(1)}\left[y(k)-\sum_{i=1}^{k-1} h(i) x(k+1-i)\right] \tag{7-33}
\end{equation*}
$$

For $e_{s}(t){\underset{R}{g}}$, i.e., (7-54)

$$
\begin{equation*}
x(1)=1+\bar{R}_{g} \tag{7-56}
\end{equation*}
$$

while for $e_{s}(t)$, (7-38)

$$
x(1)=2
$$

Because $x(1)$ occurs in (7-33) only as the factor $1 / x(1)$ external to the brackets, the two series expansions are related by a simple scale factor,

$$
\begin{aligned}
\left.e_{s}(t)\right]_{g} & =\frac{2}{1+\bar{R}_{g}} e_{s}(t) \\
& =\frac{2 R_{o}}{R_{0}+R_{g}} e_{s}(t)
\end{aligned}
$$

which is the result given above in (7-51).
An experimental run on Cable $I$ using a two channel 50 ohm feed-through sampling system produced the data shown in fig. 7-8. Note that the initial step is given by

$$
\begin{equation*}
\left.e_{s}(0)\right]=\frac{2 \times 124}{100}\left(\frac{1}{2}\right)=0.55 \tag{7-57}
\end{equation*}
$$

where $R_{g}$ and $R_{o}$ are 100 and 124 ohms, respectively. The measurement system is shown in fig. 7-9. When the voltage of one of the balanced generators was reduced to zero with the 50 source impedance being maintained, no significant change (none was discerned) in the energized channel's sending-end voltage, $\left.(1 / 2) e_{S}(t)\right]$.

124


Figure 7-1. The RT diagram for a uniform transmission line connecting an arbitrary generator and load.


Figure 7.2 The initial received voltage $\begin{gathered}{\left[e_{r}(t)\right] \text { assuming } e_{g}(t) \text { to be the }} \\ 0<t<3 \ell T\end{gathered}$ unit-step. Equation (7-21) is the Laplace transform of this initial voltage. Note the transmission delay of $\ell T$ seconds and then the emergence of the distorted step-like signal from the transmission line, $\ell T$ to $3 \ell T$ seconds.


Figure 7-3. The initial sending-end voltage $\left[e_{s}(t)\right]$ assuming $e_{g}(t)$ to be $0<t<2 \ell T$
the initial step. The initial voltage jumps to 0.5 volts and then rises slowly until the first reflection returns from the load.


Figure 7-4A. The voltage divider circuit representing the voltage transfer from $e_{g}(t)$ to $e_{s}(t)$ for $0<t<2 l T$ or for all $t$ when $\ell=\infty$, or the line is terminated in $Z_{o}(s)$.


Figure 7-4B. The sending-end voltage step response when the line is infinitely long or terminated in $Z_{o}(s)$.
1.0 v
0.5 v
0.0 v
s7T0

0.5 V
0.0
$877^{\circ} \Lambda$


Time, microseconds

Figure 7-7. The divergent properties of the series (7-38) using the parameters for Cable I. The parametric variable 10 corresponds to 10 terms comprised of $0.5 \mathrm{u}(\mathrm{t})$ plus the algebraic sum of the weighted values of $w_{1}(t)$ through $w_{3}(t)$. Similarilarly, 15 corresponds to 15 terms obtained from the sum through $w_{4}(t)$; and so on through $w_{7}(t)$ for 36 terms.
1.0 V
0.5 V
0.0
s7ton

$50 \Omega$

Figure 7-9 The measurement system for observing the sending-end step response. $e_{g}(t)$ is a unit step generator.

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## A. 1 Cable Parameter Tables


Table $A-1$., continued

| Cable | Signal Conductor(s) Material and O.D. | Dielectric <br> Material and O.D. | $\begin{gathered} \text { Shield(s) } \\ \text { Material and 0.D. } \end{gathered}$ | $\begin{gathered} \text { Jacket } \\ \text { Material and 0.D. } \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: |
| F | stranded bare copper 0.81 mm $(0.032 \mathrm{in})$ | polyethylene 4.11 mm (0.162 in) | tinned copper braid 4.70 mm (0.185 in) | $\begin{gathered} \text { black pvc } \\ 6.22 \mathrm{~mm} \\ (0.245 \mathrm{in}) \end{gathered}$ |
| G | stranded bare copper 1.35 mm (0.053 in) | polyethylene <br> 7.21 mm (0.284 in) | tinned copper braid 8.08 mm (0.318 in) | $\begin{gathered} \text { black pvc } \\ 10.67 \mathrm{~mm} \\ (0.420 \mathrm{in}) \end{gathered}$ |
| H | stranded bare copper 0.56 mm (0.022 in) | $\begin{gathered} \text { polyethylene } \\ 4.11 \mathrm{~mm} \\ (0.162 \mathrm{in}) \end{gathered}$ | tinned copper <br> 4.70 mm (0.185 in) | $\begin{gathered} \text { black pvc } \\ 6.22 \mathrm{~mm} \\ (0.245 \mathrm{in}) \end{gathered}$ |
| I | stranded bare copper 0.97 mm (0.038 in) | $\begin{gathered} \text { polyethylene } \\ 7.21 \mathrm{~mm} \\ (0.284 \mathrm{in}) \end{gathered}$ | tinned copper braid 8.08 mm (0.318 in) | black pvc 10.67 mm $(0.420 \mathrm{in})$ |
| J | stranded tinned copper 0.56 mm (0.022 in) | polyethylene 2.06 mm (0.081 in) | bare copper braid 4.95 mm (0.195 in) | $\begin{gathered} \text { blue pvc } \\ 6.22 \mathrm{~mm} \\ (0.245 \mathrm{in}) \end{gathered}$ |
| K | stranded bare copper 2.74 mm (0.108 in) | foam polyethylene $(0.285 \mathrm{~mm})$ | bare copper braid (2) | $\begin{gathered} \text { black polyethylene } \\ 12.19 \mathrm{~mm} \\ (0.480 \mathrm{in}) \end{gathered}$ |
| WD-37 | stranded bare copper 0.762 mm (0.030 in) | $\begin{gathered} \text { polypropylene } \\ 2.29 \mathrm{~mm} \\ (0.09 \mathrm{in}) \end{gathered}$ | bare copper braid 2.97 mm (0.117 in) | black pvc |

Table A-2. Cable Electrical Parameters

| Cable | Capacitance per meter (ft) in pf | Inductance per meter (ft) in $n h$ | Resistance per meter (ft) in $m \Omega$ | $\begin{aligned} & \text { Loss } \\ & \text { slope } \\ & \text { "m" } \end{aligned}$ | $\begin{gathered} \text { Loss } \\ \text { constant } \\ " K "\left(x 10^{-4}\right) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $R G-58 C / U$ | 101.0 (30.8) | 252.6 (77.0) | 37.1 (11.3) | 0.52975 | 0.16710 |
| RG-214/U | 101.0 (30.8) | 252.6 (77.0) | 10.3 (3.15) | 0.55071 | 0.053034 |
| $R G-223 / U$ | 101.0 (30.8) | 252.6 (77.0) | 34.8 (10.6) | 0.52314 | 0.16676 |
| $R G-59 B / U$ | 67.6 (20.6) | 380.2 (115.9) | 147.6 (45.0) | 0.52284 | 0.21993 |
| A | 64.3 (19.6) | 391.1 (119.2) | 58.7 (17.9) | 0.50981 | 0.62760 |
| B | 64.3 (19.6) | 391.1 (119.2) | 20.0 (6.1) | 0.52724 | 0.22117 |
| C | 64.6 (19.7) | 393.4 (119.9) | 68.2 (20.8) | 0.51395 | 0.71262 |
| D | 64.6 (19.7) | 393.4 (119.9) | 196.9 (60.0) | 0.55307 | 0.46238 |
| E | 68.9 (21.0) | 419.0 (127.7) | 68.9 (21.0) | 0.52174 | 0.51676 |
| RG-22B/U | 52.5 (16.0) | 473.8 (144.4) | 43.0 (13.1) | 0.53103 | 0.23926 |
| F | 51.2 (15.6) | $491.5(149.8)$ | 86.6 (26.4) | 0.51565 | 0.51543 |
| G | 51.2 (15.6) | 491.5 (149.8) | 31.5 (9.6) | 0.54819 | 0.17271 |
| H | 40.4 (12.3) | 620.4 (189.1) | 160.1 (48.8) | 0.50262 | 0.78542 |
| I | 40.4 (12.3) | 620.4 (189.1) | 61.7 (18.8) | 0.53952 | 0.25639 |
| J | 40.4 (12.3) | 620.4 (189.1) | 186.4 (56.8) | 0.51619 | 0.69930 |
| K | 85.3 (26.0) | 213.3 (65.0) | 7.45 (2.27) | 0.55829 | 0.039593 |
| WD-37 | 43.6 (13.3) | $587.2(179.0)$ | $137.8(42.0)$ | 0.57232 | 0.086936 |



Figure A-1. Modeled attenuation plot for 304.8 meters ( 1000 feet) of RG-58C/U. Ordinate units are decibels and abscissa units are megahertz.



Figure A-3. Modeled and computed time domain impulse response for 304.8 meters ( 1000 feet) of RG-58C/U. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .
$1.09 \times 10^{-2}$


Figure A-4. Time-expanded modeled and computed time domain impulse response for 304.8 meters ( 1000 feet) of RG-58C/U. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns.


Figure A-5. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of $\mathrm{RG}-58 \mathrm{C} / \mathrm{U}$. Ordinate units are volts and abscissa units are nanoseconds.


Figure A-6. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable RG-58C/U. Ordinate units are volts and abscissa units are hertz.


Figure A-7. TDR unit step response for cable RG-58C/U. Source step generator has assumed 50 resistive output impedance. Ordinate units are volts and abscissa units are microseconde.


Figure A-8. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of RG-214/U. Ordinate units are decibels and abscissa units are megahertz.


Figure A-9. Modeled minimum-phase phase shift plot for 304.8 meters ( 1000 ft ) of RG-214/U. Ordinate units are degrees and abscissa units are megahertz.


[^1]$2.46 \times 10^{-2}$


Figure A-11. Time-expanded modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of RG-214/U. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A-12. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of RG-214/U. Ordinate units are volts and abscissa units are nanoseconds.


Figure A-13. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable $\mathrm{RG}-214 / \mathrm{U}$. Ordinate units are volts and abscissa units are


Figure A-14. TDR unit step response for cable RG-214/U. Source step generator has assumed $50 \Omega$ resistive output impedance. Ordinate units are volts and abscissa units are microseconds.


Figure A-15. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of RG-223/U. Ordinate units are decibels and abscissa units are megahertz.


Figure A-16. Modeled minimum-phase phase shift plot for 304.8 meters ( 1000 ft ) of RG-223/U. Ordinate units are degrees and abscissa units are megahertz.
$1.35 \times 10^{-2}$
Figure A-17. Modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of
RG-223/U. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing
between points is 0.9766 ns.
Figure A-18. Time-expanded modeled and computed time domain impulse response for 304.8 meters ( 1000
ft) of RG-223/U Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time
spacing between points is 0.9766 ns.


Figure A-19. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of RG-223/U. Ordinate units are volts and abscissa units are nanoseconds.


Figure A-20. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable RG-223/U. Ordinate units are volts and abscissa units are hertz.

# 100 <br>  <br> Figure A-21. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of $\mathrm{RG}-59 \mathrm{~B} / \mathrm{U}$. Ordinate units are decibels and abscissa units are megahertz. 



Figure A-22. Modeled minimum-phase phase shift plot for 304.8 meters ( 1000 ft ) of RG-59B/U. Ordinate units are degrees and abscissa units are megahertz.


Figure A-23. Modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of RG-59B/U. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .
$2.18 \times 10^{-2}$


Figure A-24. Time-expanded modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of $\mathrm{RG}-59 \mathrm{~B} / \mathrm{U}$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A-25. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of RG-59B/U. Ordinate units are volts and abscissa units are nanoseconds.


Figure A-26. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable $\mathrm{RG}-59 \mathrm{~B} / \mathrm{U}$. Ordinate units are volts and abscissa units are hertz.


Figure A-27. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of A. Ordinate units are decibels and abscissa units are megahertz.


Figure A-28. Modeled minimum-phase phase shift plot for 304.8 meters ( 1000 ft ) of A. Ordinate units are degrees and abscissa units are megahertz.


Figure A -29. Modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of A . Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A-30. Time-expanded modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of $A$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A-31. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of A . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-32. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable $A$. Ordinate units are volts and abscissa units are hertz.


Figure A-33. Modeled attenuation plot for 60.96 meters ( 200 ft ) of B. Ordinate units are decibels and abscissa units are megahertz.


Figure A-34. Modeled minimum-phase phase shift plot for 60.96 meters ( 200 ft ) of B. Ordinate units are degrees and abscissa units are megahertz.

Figure A-35. Modeled and computed time domain impulse response for 60.96 meters (200 ft) of $B$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between
(109 $100^{-2}$
Figure A-36. Time-expanded modeled and computed time domain impulse response for 60.96 meters (200 ft) of $B$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure $A$-37. Modeled and computed time domain unit step response for 60.96 meters (200 ft ) of B . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-38. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 60.96 meters ( 200 ft ) of cable B. Ordinate units are volts and abscissa units are hertz.


Figure A-39. Modeled attenuation plot for 152.4 meters ( 500 ft ) of B. Ordinate units are decibels and abscissa units are megahertz.


Figure A-40. Modeled minimum-phase phase shift plot for 152.4 meters ( 500 ft ) of B. Ordinate units are degrees and abscissa units are megahertz.


$$
\begin{aligned}
& 1.147 \times 10^{-2} \\
& \text { Figure A-42. Time-expanded modeled and computed time domain impulse response for } 152.4 \text { meters ( } 500 \mathrm{ft} \text { ) } \\
& \text { of } B \text {. Ordinate units are seconds }{ }^{-1} \text {, abscissa units are nanoseconds and time spacing } \\
& \text { between points is } 0.9766 \text { ns. }
\end{aligned}
$$

## 

Figure A-43. Modeled and computed time domain unit step response for 152.4 meters ( 500 ft ) of B . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-44. Plots of "zero"/"one" cable unit step response voltages versus log 10 frequency for 152.4 meters ( 500 ft ) of cable B. Ordinate units are volts and abscissa units are hertz.


Figure $A-45$. Modeled attenuation plot for 333.76 meters ( 1095 ft ) of $B$. Ordinate units are decibels and abscissa units are megahertz.


Figure A-46. Modeled minimum-phase phase shift plot for 333.76 meters ( 1095 ft ) of B. Ordinate units are degrees and abscissa units are megahertz.



Figure A-48. Time-expanded modeled and computed time domain impulse response for 333.76 meters (l095 $f t$ ) of $B$. Ordinate units are seconds ${ }^{-l}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A-49. Modeled and computed time domain unit step response for 333.76 meters ( 1095 ft ) of B . ordinate units are volts and abscissa units are nanoseconds.


Figure A-50. Plots of "zero"/"one" cable unit step response voltages versus loglo frequency for 333.76 meters ( 1095 ft ) of cable B. Ordinate units are volts and abscissa units are hertz.


Figure A-51. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of C. Ordinate units are decibels and abscissa units are megahertz.

$1.37 \times 10^{-2}$

Figure A-53. Modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of C . Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .

Figure A-54. Time-expanded modeled and computed time domain impulse response for 304.8 meters (1000 $f t$ ) of $C$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns.


Figure $A-55$. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of C . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-56. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable $C$. Ordinate units are volts and abscissa units are hertz.


Figure A-57. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of D. Ordinate units are decibels and abscissa units are megahertz.


Figure A-58. Modeled minimum-phase phase shift plot for 304.8 meters ( 1000 ft ) of D. Ordinate units are degrees and abscissa units are megahertz.
 Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 4.883 ns.
$8.42 \times 10^{-3}$


Figure A-60. Time-expanded modeled and coinputed time domain impulse response for 304.8 meters ( 1000 $f t$ ) of $D$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 4.883 ns.


Figure A-61. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of D . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-62. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable D. Ordinate units are volts and abscissa units are hertz.


Figure A-63. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of E. Ordinate units are decibels and abscissa units are megahertz.


Figure A-64. Modeled minimum-phase phase shift plot for 304.8 meters ( 1000 ft ) of E. Ordinate units are degrees and abscissa units are megahertz.


Figure A-65. Modeled and computed time domain impulse response for 304.8 meters (1000 ft) of E. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 4.883 ns.


Figure A-66. Time-expanded modeled and computed time domain impulse response for 304.8 meters (1000 $f t$ ) of $E$ Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 4.883 ns .


Figure A-67. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of E . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-68. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable E . Ordinate units are volts and abscissa units are hertz.


Figure $A-70$. Modeled minimum-phase phase shift plot for 60.96 meters (200 ft) of RG-22B/U. Ordinate units are degrees and abscissa units are megahertz.


Figure A-71. Modeled and computed time domain impulse response for 60.96 meters ( 200 ft ) of RG-22B/U. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns.


Figure A-72. Time-expanded modeled and computed time domain impulse response for 60.96 meters (200 ft) of $R G-22 B / U$ Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns.


Figure A-73. Modeled and computed time domain unit step response for 60.96 meters ( 200 ft ) of RG-22B/U. Ordinate units are volts and abscissa units are nanoseconds.


Figure A-74. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 60.96 meters ( 200 ft ) of cable $\mathrm{RG}-22 \mathrm{~B} / \mathrm{U}$. Ordinate units are volts and abscissa units are hertz.

# 1 <br> $\begin{array}{cc}1 & 1 \\ 1 & 10\end{array}$ <br> 100 <br> 1000 <br> 10000 

Figure A-75. Modeled attenuation plot for 152.4 meters ( 500 ft ) of RG-22B/U. Ordinate units are decibels and abscissa units are megahertz.


Figure A-76. Modeled minimum-phase phase shift plot for 152.4 meters ( 500 ft ) of RG-22B/U. Ordinate units are degrees and abscissa units are megahertz.
$1.367 \times 10^{-2}$
Figure A-77. Modeled and computed time domain impulse response for 152.4 meters ( 500 ft ) of
RG-22B/U. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing
between points is 0.1953 ns .

$$
\begin{aligned}
& 1.367 \times 10^{-2} 1 \\
& \text { Figure A-78. Time-expanded modeled and computed time domain impulse response for } 152.4 \text { meters ( } 500 \mathrm{ft} \text { ) } \\
& \text { of } R G-22 B / U \text {. Ordinate units are seconds }{ }^{-1} \text {, abscissa units are nanoseconds and time } \\
& \text { spacing between points is } 0.1953 \mathrm{~ns} \text {. }
\end{aligned}
$$



Figure A-79. Modeled and computed time domain unit step response for 152.4 meters ( 500 ft ) of RG-22B/U. Ordinate units are volts and abscissa units are nanoseconds.


Figure A-80. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 152.4 meters ( 500 ft ) of cable RG-22B/U. Ordinate units are volts and abscissa units are hertz.


Figure A-81. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of $\mathrm{RG}-22 \mathrm{~B} / \mathrm{U}$. Ordinate units are decibels and abscissa units are megahertz.


Figure A-82. Modeled minimum-phase phase shift plot for 304.8 meters ( 1000 ft ) of RG-22B/U. Ordinate units are degrees and abscissa units are megahertz.


Figure A-83. Modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of RG-22B/U. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A-84. Time-expanded modeled and computed time domain impulse response for 304.8 meters ( 1000 $f t$ ) of $R G-22 B / U$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns.

> Figure A-85. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of RG$22 \mathrm{~B} / \mathrm{U}$. Ordinate units are volts and abscissa units are nanoseconds.


Figure A-86. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable RG-22B/U. Ordinate units are volts and abscissa units are hertz.


Figure $A-87$. TDR unit step response for cable $R G-22 B / U$. Source step generator has assumed $50 \Omega$ resistive output impedance. Ordinate units are volts and abscissa units are microseconds.


Figure A-88. Modeled attenuation plot for 320.04 meters ( 1050 ft ) of F. Ordinate units are decibels and abscissa units are megahertz.


Figure A-89. Modeled minimum-phase phase shift plot for 320.04 meters ( 1050 ft ) of F . Ordinate units are degrees and abscissa units are megahertz.


Figure A -90. Modeled and computed time domain impulse response for 320.04 meters ( 1050 ft ) of F . Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 1.953 ns .


Figure A-91. Time-expanded modeled and computed time domain impulse response for 320.04 meters ( 1050 $f t$ ) of $F$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 1.953 ns .


Figure A-92. Modeled and computed time domain unit step response for 320.04 meters ( 1050 ft ) of F . Ordinate units are volts and abscissa units are nanoseconds.

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Figure A-93. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 320.04 meters ( 1050 ft ) of cable $F$. Ordinate units are volts and abscissa units are hertz.


Figure A-94. Modeled attenuation plot for 323.09 meters ( 1060 ft ) of G. Ordinate units are decibels and abscissa units are megahertz.


Figure A-95. Modeled minimum-phase phase shift plot for 323.09 meters ( 1060 ft ) of G. Ordinate units are degrees and abscissa units are megahertz.



Figure A-97, Time-expanded modeled and computed time domain impulse response for 323.09 meters (1060 ft ) of $G$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A-98. Modeled and computed time domain unit step response for 323.09 meters ( 1060 ft ) of G . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-99. Plots of "zero"/" one" cable unit step response voltages versus $\log _{10}$ frequency for 323.09 meters ( 1060 ft ) of cable $G$. Ordinate units are volts and abscissa units are hertz.

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Figure $A^{-}$100. Modeled attenuation plot for 152.4 meters ( 500 ft ) of H . Ordinate units are decibels and abscissa units are megahertz.


Figure A- 101. Modeled minimum-phase phase shift plot for 152.4 meters ( 500 ft ) of H . Ordinate units are degrees and abscissa units are megahertz.


Figure A-102. Modeled and computed time domain impulse response for 152.4 meters ( 500 ft ) of H . Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 4.883 ns.


Figure A-103. Time-expanded modeled and computed time domain impulse response for 152.4 meters (500 ft) of $H$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 4.883 ns .


Figure A-104. Modeled and computed time domain unit step response for 152.4 meters ( 500 ft ) of H . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-105. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 152.4 meters ( 500 ft ) of cable H . Ordinate units are volts and abscissa units are hertz.


Figure A-106. Modeled attenuation plot for 326.14 meters ( 1070 ft ) of H. Ordinate units are decibels and abscissa units are megahertz.


Figure A-107. Modeled minimum-phase phase shift plot for 326.14 meters ( 1070 ft ) of H. Ordinate units are degrees and abscissa units are megahertz.



Figure A-109. Time-expanded modeled and computed time domain impulse response for 326.14 meters ( 1070 $f t$ ) of H. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 4.883 ns .


Figure A-110. Modeled and computed time domain unit step response for 326.14 meters ( 1070 ft ) of H . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-111. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 326.14 meters ( 1070 ft ) of cable H . Ordinate units are volts and abscissa units are hertz.


Figure A-112. TDR unit step response for cable $H$. Source step generator has assumed $50 \Omega$ resistive output impedance. Ordinate units are volts and abscissa are microseconds.


Figure A-113. Modeled attenuation plot for 60.96 meters ( 200 ft ) of I. Ordinate units are decibels and abscissa units are megahertz.


Figure A- 114 Modeled minimum-phase phase shift plot for 60.96 meters ( 200 ft ) of I. Ordinate units are degrees and abscissa units are megahertz.


Figure $A-115$. Modeled and computed time domain impulse response for 60.96 meters (200 ft ) of I . ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A- 116. Time-expanded modeled and computed time domain impulse response for 60.96 meters ( 200 ft ) of I. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns.


Figure A- 117. Modeled and computed time domain unit step response for 60.96 meters ( 200 ft ) of I . Ordinate units are volts and abscissa units are nanoseconds.

> Figure A- 118. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 60.96 meters (200 ft) of cable I. Ordinate units are volts and abscissa units are hertz.


Figure A-119. Modeled attenuation plot for 152.4 meters ( 500 ft ) of . Ordinate units are decibels and abscissa units are megahertz.

Figure A-120. Modeled minimum-phase phase shift plot for 152.4 meters ( 500 ft ) of I. Ordinate units are degrees and abscissa units are megahertz.



Figure A-122. Time-expanded modeled and computed time domain impulse response for 152.4 meters (500 ft) of I. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A-123. Modeled and computed time domain unit step response for 152.4 meters ( 500 ft ) of I . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-124. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 152.4 meters ( 500 ft ) of cable I. Ordinate units are volts and abscissa units are hertz.


Figure A-125. Modeled attenuation plot for 320.04 meters ( 1050 ft ) of I. Ordinate units are decibels and abscissa units are megahertz.


Figure A-126. Modeled minimum-phase phase shift plot for 320.04 meters ( 1050 ft ) of I. Ordinate units are degrees and abscissa units are megahertz.


Figure A-127. Modeled and computed time domain impulse response for 320.04 meters ( 1050 ft ) of I . ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .

\section*{$1.831 \times 10^{-2}$ <br> 

Figure A-128. Time-expanded modeled and computed time domain impulse response for 320.04 meters ( 1050 $f t$ ) of $I$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .


Figure A-129. Modeled and computed time domain unit step response for 320.04 meters ( 1050 ft ) of I . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-130. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 320.04 meters ( 1050 ft ) of cable I . Ordinate units are volts and abscissa units are hertz.


Figure A-131. TDR unit step response for cable I. Source step generator has assumed $50 \Omega$ resistive output impedance. Ordinate units are volts and abscissa units are microseconds.


Figure A-132. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of J. Ordinate units are decibels and abscissa units are megahertz.


Figure A-133. Modeled minimum-phase phase shift plot for 304.8 meters ( 1000 ft ) of J. Ordinate units are degrees and abscissa units are megahertz.

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\begin{aligned}
& \text { 2. } 80 \times 10^{-2} \\
& \text { Figure A-134. Modeled and computed time domain impulse response for } 304.8 \text { meters ( } 1000 \mathrm{ft} \text { ) of } \mathrm{J} \text {. } \\
& \text { Ordinate units are seconds }{ }^{-1} \text {, abscissa units are nanoseconds and time spacing between } \\
& \text { 2. } 80 \times 10^{-2} \\
& \text { Figure A-135. Time-expanded modeled and computed time domain impulse response for } 304.8 \text { meters ( } 1000 \\
& f t \text { ) of J. Ordinate units are seconds }{ }^{-1} \text {, abscissa units are nanoseconds and time spacing } \\
& \text { between points is } 0.9766 \mathrm{~ns} \text {. }
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Figure $A-137$. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable J. Ordinate units are volts and abscissa units are hertz.

Figure A-138. Modeled attenuation plot for 304.8 meters ( 1000 ft ) of K. Ordinate units are decibels and abscissa units are megahertz.

Figure A-139. Modeled minimum-phase phase shift plot for 304.8 meters ( 1000 ft ) of K . Ordinate units are degrees and abscissa units are megahertz.


Figure A-140. Modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of K . Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.4883 ns .


Figure A-141. Time-expanded modeled and computed time domain impulse response for $304: 8$ meters (l000 ft ) of K . Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .

Figure A-142. Modeled and computed time domain unit step response for 304.8 meters ( 1000 ft ) of K . Ordinate units are volts and abscissa units are nanoseconds.


Figure A-143. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 304.8 meters ( 1000 ft ) of cable K . Ordinate units are volts and abscissa units are hertz.


Figure A-144. Modeled attenuation plot for 402.34 meters ( 1320 ft ) of $\mathrm{WD}-37$. Ordinate units are decibels and abscissa units are megahertz.


Figure A-145. Modeled minimum-phase phase shift plot for 402.34 meters ( 1320 ft ) of WD-37. Ordinate units are degrees and abscissa units are megahertz.


Figure A-146. Modeled and computed time domain impulse response for 402.34 meters ( 1320 ft ) of $\mathrm{WD}-37$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 1.953 ns .


Figure A-147. Time-expanded modeled and computed time domain impulse response for 402.34 meters (1320 ft ) of $\mathrm{WD}-37$. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 1.953 ns.


Figure A-148. Modeled and computed time domain unit step response for 402.34 meters ( 1320 ft ) of WD 37. Ordinate units are volts and abscissa units are nanoseconds.


Figure A-149. Plots of "zero"/"one" cable unit step response voltages versus $\log _{10}$ frequency for 402.34 meters ( 1320 ft ) of $\mathrm{WD}-37$. Ordinate units are volts and abscissa units are hertz.

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Figure $A-150$. TDR unit step response for cable $W D-37$. Source step generator has assumed 50 resistive output impedance. Ordinate units are volts and abscissa units are microseconds.
B.1. APMS System Calibration Programs and Subroutines
B.1.1. Y-Axis Calibration Subroutine






THIS SUBROUTINE ACQUIRES Y-AXIS OATA FROM THE SAMPLINE
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FRON THE MAIN PROGRAM BY THE FORTRAN STATEMENT,

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B.l.5. General Matrix Multiplication Subroutine


B.2. Cable Measurement Programs and Subroutines
B.2.1. Main Cable Time Domain Transfer Function Measurement Program


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B.2.5. General Time Domain/Frequency Domain Plot Program




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(IT.EQ.4) ACCEPT MRADIANS (1) OR DEGEES (2)TM.LEZ
(LGI.EQ.2) GO TO 210
60 TO 220
co TO 225
co
Q.2)
3-GRINF. 4-rMr.

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B.3. Cable Modeling and Analysis Programs
B.3.1. Main Cable Model Program for Calculations of "m", "K", and Cable Impulse Response
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B.3.2. Main Cable Model Program for Calculation of Cable Step Response, Bit Error Waveform and Cable Square Wave Response

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B.3.4. Main Program for Calculation of Cable Matching Network Resistors


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B.3.6. Function Subroutine for Calculation of the Gamma Function

B.3.7. Subroutine for Accurate Calculations of Cable Time Domain Impulse Response dc Level


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Figure $C .1$ is a simplified block diagram of the Automatic Pulse Measurement System (APMS) [6,7]. The sampling oscilloscope is a commercially available unit that has been modified to allow either stand-alone or computer-controlled operation. It should be noted that this type of oscilloscope is not a real time instrument. Rather, it acquires one voltage versus time point of the measured waveform per waveform occurrence. Thus only repetitive waveforms may be observed. The advantage of this sampling scheme over conventional real time oscilloscopes is the large amount of available bandwidth. Whereas conventional oscilloscopes are presently limited in bandwidth to below 1 GHz , the effective bandwidth of the sampling cscilloscope used in the AFMS is about 18 GHz .

When the sampling oscilloscope is switched to computer-controlled operation, the oscilloscope is connected to the system minicomputer via an $A / D-D / A$ unit and a digital control interface. Under program control, the computer sends the oscilloscope the desired time location at which a voltage sample is to be taken via the 14 -bit D/A converter. After the waveform voltage sample has been taken it is recorded in the computer memory via the 14 -bit $A / D$ converter. The digital interface is designed to allow asynchronous (independent) clocking of the sampling oscilloscope with respect to the minicomputer.

The system minicomputer is a 16-bit machine with 32 K words of core memory. With the peripherals shown in the block diagram, i.e., the flexible disk drive, the CRT graphics terminal, hard copy unit, paper tape punch (PTP) and paper tape reader (PTR), the system can generate either alpha-numeric or graphical data in either soft or hard copy. In addition, a hardware floating point number processor has been added to the computer to speed arithmetic calculations by an approximate factor of 10 .

The APMS is capable of running at a maximum sampling rate of approximately 8 kHz . The useful signal voltage measurement range is $\pm 800 \mathrm{mV}$ into the $50 \Omega$ sampling head. The available real time equivalent sweep speeds range from $10 \mathrm{ps} / \mathrm{cm}$ to $500 \mu \mathrm{~s} / \mathrm{cm}$ in a 1-2-5 sequence.

The system software consists of a commercially available flexible disk operating system with such features as Basic and Fortran language support, disk files, text editor, and macro-command capability. The sampling oscilloscope waveform acquisition is done with assembly language programming while virtually all signal processing of the waveform is done in either Basic or Fortran. The software has been designed to achieve a happy medium between system versatility and ease of operation.


Figure C. 1 Simplified block diagram of APMS

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[^1]:    Figure A-10. Modeled and computed time domain impulse response for 304.8 meters ( 1000 ft ) of RG-214/U. Ordinate units are seconds ${ }^{-1}$, abscissa units are nanoseconds and time spacing between points is 0.9766 ns .

